

Problem set 4 for Quantum Field Theory course

2019.03.05.

Topics covered

- Representation of C, P, T transformations on Dirac field and bilinears
- Gordon identity and current decomposition

Recommended reading

Peskin–Schroeder: An introduction to quantum field theory

- Section 3.6

Problem 4.1 P transformation of Dirac field

Recall that the Dirac field can be expressed as

$$\Psi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_{\mathbf{p}}}} \sum_{s=1}^2 \left[a_{\mathbf{p}}^s u^s(\mathbf{p}) e^{-ipx} + b_{\mathbf{p}}^{s\dagger} v^s(\mathbf{p}) e^{ipx} \right], \quad (1)$$

with two particle species a and b . Under parity, annihilation operators transform as

$$Pa_{\mathbf{p}}^s P^{-1} = \eta_a a_{-\mathbf{p}}^s, \quad Pb_{\mathbf{p}}^s P^{-1} = \eta_b b_{-\mathbf{p}}^s. \quad (2)$$

Show that the action of parity on the spinor field $\Psi(t, \mathbf{x})$ can be written as

$$P\Psi(t, \mathbf{x})P^{-1} = \mathcal{P}\Psi(t, -\mathbf{x}). \quad (3)$$

where the constant matrix \mathcal{P} is given by

$$\mathcal{P} = \eta_a \gamma^0. \quad (4)$$

(Hint: you may find useful the relation $\tilde{p}\sigma = p\bar{\sigma}$, where $\tilde{p}^\mu = (p^0, -\mathbf{p})$.)

Note that locality of the parity transformed field leads to a relation between η_a and η_b ! What is this relation?

From Eq. (3) derive the analogous transformation rule for $\bar{\Psi}$!

Problem 4.2 T transformation of Dirac field

Both momentum and spin must change sign under time reversal, therefore we can write

$$Ta_{\mathbf{p}}^s T^{-1} = a_{-\mathbf{p}}^{-s}, \quad Tb_{\mathbf{p}}^s T^{-1} = b_{-\mathbf{p}}^{-s}, \quad (5)$$

where $-s$ refers to the flipped spin.

Remark: in principle one could also allow for a phase ζ in this definition. However, as it was shown in the previous problem set, it is physically irrelevant. A possible choice is such that product of phases in CPT combination is equal to one.

- (a) We have to implement the spin flip on spinors. Prove the identity

$$\vec{\sigma}\sigma^2 = \sigma^2(-\vec{\sigma})^*, \quad (6)$$

where $\vec{\sigma} = (\sigma^1, \sigma^2, \sigma^3)^T$ is a vector of Pauli matrices.

- (b) Consider a spinor ξ that has +1 spin projected on axis \vec{n} :

$$\vec{n}\vec{\sigma}\xi = \xi. \quad (7)$$

Utilize the identity (6) to show that operator $-i\sigma^2$ flips the spin (i.e. the eigenvalue of ξ under $\vec{n}\vec{\sigma}$), so

$$\xi^{-s} = -i\sigma^2\xi^{s*}. \quad (8)$$

- (c) Using Eq. (6) show that

$$u^{-s}(\tilde{p}) = -\gamma^1\gamma^3u^s(p)^*. \quad (9)$$

(Hint: think about the projector resolution of $\sqrt{p\sigma}$.)

- (d) After all this preparation, show using (1) that

$$T\Psi(t, \mathbf{x})T^{-1} = -\gamma^1\gamma^3\Psi(-t, \mathbf{x}). \quad (10)$$

Problem 4.3 C transformation of Dirac field

Charge conjugation acts on operators as follows

$$Ca_{\mathbf{p}}^s C^{-1} = \chi_a b_{\mathbf{p}}^s, \quad Cb_{\mathbf{p}}^s C^{-1} = \chi_b a_{\mathbf{p}}^s, \quad (11)$$

so it relates the two species, i.e. particle and anti-particle.

- (a) Using again Eq. (6) prove that positive and negative energy spinors are connected via the relation

$$u^s(\mathbf{p}) = -i\gamma^2(v^s(\mathbf{p}))^*, \quad (12)$$

where in $u^s(p)$ we have the 2-spinor ξ^s while in $v^s(p)$ there is ξ^{-s} in accordance with what we learned about spin projections in Problem 2.5 e).

- (b) Using (1), derive the action of this operator on the Dirac spinor field

$$C\Psi C^{-1} \equiv \Psi^c = \chi_a C\bar{\Psi}^T. \quad (13)$$

Note that locality of the charge conjugate field leads to a relation between χ_a and χ_b ! What is this relation?

Show that in Dirac representation

$$C = -i\gamma^2\gamma^0 \quad (14)$$

i.e. identical (up to a sign) to the charge conjugation matrix obtained in Problem 2.6.

Perform a similar calculation (or use Eq. (13)) and derive the expression

$$C\bar{\Psi}C^{-1} \equiv \bar{\Psi}^c = -\chi_a^*\Psi^T C. \quad (15)$$

Problem 4.4 C, P, T transformation of Dirac field bilinears

- (a) The transformation properties of Dirac spinors under discrete symmetry transformations were derived in Problem 4.1. Combine the results to obtain the action of CPT on a Dirac spinor.
- (b) Show that the product $\bar{\Psi}\Psi$ transforms as a scalar under each discrete symmetry.
- (c) Derive the transformation properties of $\bar{\Psi}\gamma^\mu\Psi$, $\bar{\Psi}\sigma^{\mu\nu}\Psi$, $\bar{\Psi}\gamma^\mu\gamma^5\Psi$ and $\bar{\Psi}\gamma^5\Psi$ under C , P and T . Note that the P and T transformation show a particular dependence on Lorentz indices.

- (d) Obtain the CPT transformation of the Dirac bilinears.
(e) How does ∂_μ transform under C , P and T ?
(f) Now we are ready to compute the effect of CPT on the free Dirac Lagrangian \mathcal{L}_0

$$\mathcal{L}_0(x) = \bar{\Psi}(x)(i\gamma^\mu\partial_\mu - m)\Psi(x). \quad (16)$$

Is the action invariant under CPT ?

Problem 4.5 Gordon identity

In this exercise we prove a useful identity and then consider some of its applications.

- (a) Prove the Gordon identity

$$\bar{u}(p')\gamma^\mu u(p) = \bar{u}(p')\frac{(p+p')^\mu + i\sigma^{\mu\nu}(p'-p)_\nu}{2m}u(p), \quad (17)$$

where $u(p)$ are positive energy Dirac spinors.

Hint: starting from the right hand side rewrite $\sigma^{\mu\nu}$ using the Clifford algebra and exploit the Dirac equation in momentum space.

- (b) Recall that the $U(1)$ Noether current of Dirac field is

$$j^\mu = \bar{\Psi}\gamma^\mu\Psi. \quad (18)$$

Now consider a positive energy wave packet

$$\Psi^+(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_{\mathbf{p}}}} \sum_{s=1}^2 a_{\mathbf{p}}^s u^s(\mathbf{p}) e^{-ipx}, \quad (19)$$

where the coefficients $a_{\mathbf{p}}^s$ are complex numbers, and calculate the current $j^{\mu(+)}$ associated to it. Show that if the total charge is normalised to unity

$$\int d^3x j^{0(+)}(t, \mathbf{x}) = 1 \quad (20)$$

one can express the total current as

$$\mathbf{J}^+ = \int d^3x \mathbf{j}^{(+)}(t, \mathbf{x}) = \sum_s \int d^3p |a_{\mathbf{p}}^s|^2 \frac{\mathbf{p}}{E_{\mathbf{p}}} = \left\langle \frac{\mathbf{p}}{E_{\mathbf{p}}} \right\rangle, \quad (21)$$

i.e. it equals the group velocity.

Hint: use the Gordon-identity and orthogonality of spinors (cf. Problem 2.3(f)).

- (c) Gordon decomposition of Dirac current

Decompose the current (18) into two parts as

$$j^\mu = \frac{1}{2}(\bar{\Psi}\gamma^\mu\Psi + \bar{\Psi}\gamma^\mu\Psi), \quad (22)$$

and use free Dirac equation in the first term and the conjugate equation in the second. Writing the $\gamma^\mu\gamma^\nu$ terms in the form

$$\gamma^\mu\gamma^\nu = \eta^{\mu\nu} + \frac{1}{i}\sigma^{\mu\nu} \quad (23)$$

derive the Gordon decomposition:

$$j^\mu = \frac{i}{2m}\bar{\Psi}\overleftrightarrow{\partial}^\mu\Psi + \frac{1}{2m}\partial_\nu(\bar{\Psi}\sigma^{\mu\nu}\Psi). \quad (24)$$

Note that the first term is similar to scalar $U(1)$ current and so the second can be attributed to the spin. The coupling to an external electromagnetic field corresponds to an interaction term

$$H_{int} = - \int d^3x e j^\mu A_\mu, \quad (25)$$

where e is the electric charge, so $e j^\mu$ is the electromagnetic current. Focus on the second term from Gordon decomposition and perform a partial integration to derive the expression

$$H_{int}^{(2)} = -\frac{e}{2m} \int d^3x \frac{1}{2} F_{\mu\nu} (\bar{\Psi} \sigma^{\mu\nu} \Psi) \quad (26)$$

Take positive energy solutions in the Dirac representation

$$\Psi = \begin{pmatrix} \Psi_+ \\ \frac{\vec{\sigma} \cdot \vec{p}}{p_0 + mc} \Psi_+ \end{pmatrix} \quad (27)$$

in the nonrelativistic approximation $p = |\vec{p}| \ll mc$, where the lower component is negligible compared to the upper. Show that in this limit the above interaction can be approximated at leading order in p/c as

$$H_{int}^{(2)} \simeq -\frac{e}{m} \int d^3x \vec{B} \cdot \left(\Psi_+^\dagger \frac{\vec{\sigma}}{2} \Psi_+ \right), \quad (28)$$

which is a Pauli interaction term between the spin and the external magnetic field \vec{B} . Given the non-relativistic spin operator

$$\vec{S} = \frac{\vec{\sigma}}{2}, \quad (29)$$

what is the gyromagnetic ratio?