

A simple statement about the period of periodic motions

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Let us take a one dimensional Hamiltonian, for simplicity we assume that it is time independent:

$$H(x, p) \tag{0.1}$$

We assume that the solutions are closed orbits $x(t)$ such that

$$x(t + T) = x(t), \quad \text{where} \quad T = T(E) \tag{0.2}$$

We introduce the action angle variables (φ, I) . Here

$$2\pi I = \oint dq p \tag{0.3}$$

The parameter I is time independent, because it is integrated over the full orbit. Therefore it only depends on the overall energy, which characterizes the orbits:

$$I = I(E) \tag{0.4}$$

This relation can be understood also in the other direction: the energy is constant along the path, therefore after the canonical transformation $(x, p) \rightarrow (\varphi, I)$ the Hamiltonian only depends on the action variable:

$$H(\varphi, I) = H(I) \tag{0.5}$$

The Hamilton equations are:

$$\dot{I} = 0, \quad \dot{\varphi} = \frac{dH}{dI} \tag{0.6}$$

The right hand side is a constant in time, therefore $\varphi(t) = \frac{dH}{dI}t + c$.

What is the change of φ during one period? In the lecture you learned that this is exactly 2π by definition. It follows that

$$T = \frac{2\pi}{\dot{\varphi}} = 2\pi \frac{dI}{dE} \tag{0.7}$$

Here we used again that the energy and thus H only depends on I and not on the angle. This is what we wanted to prove.

Proof of the statement that φ is changed by 2π over one period

Here we just provide this proof once again.

Let consider φ as a function of the action and the original position: $\varphi(I, q)$. We compute the change as the contour integral (integral back and forth)

$$\Delta\varphi = \oint d\varphi = \oint \frac{\partial\varphi}{\partial q} dq \tag{0.8}$$

The φ, I pair is a canonical pair, which means that

$$\frac{\partial\varphi(I, q)}{\partial q} = \frac{\partial p(I, q)}{\partial I} \quad (0.9)$$

This can be obtained from a generating function, or directly from the canonical transformation. Now we compute it directly from the canonical transformation.

The (φ, I) pair is canonical if

$$\{\varphi, I\} = \frac{\partial\varphi}{\partial q}\Big|_p \frac{\partial I}{\partial p}\Big|_q - \frac{\partial I}{\partial q}\Big|_p \frac{\partial\varphi}{\partial p}\Big|_q = 1 \quad (0.10)$$

Here both φ and I are understood as functions of (q, p) , therefore we denoted the variable which is kept fixed in the subscript. The difference with (0.9) is that in the Poisson bracket the variables are the original (q, p) , and we have to change to the mixed parametrization (q, I) .

Let us consider changing p, q such that I is kept constant. Then

$$\frac{\partial p}{\partial q}\Big|_I = -\frac{\frac{\partial I}{\partial q}\Big|_p}{\frac{\partial I}{\partial p}\Big|_q} \quad (0.11)$$

Substituting into (0.10)

$$\frac{\partial\varphi}{\partial q}\Big|_p \frac{\partial I}{\partial p}\Big|_q + \frac{\partial p}{\partial q}\Big|_I \frac{\partial I}{\partial p}\Big|_q \frac{\partial\varphi}{\partial p}\Big|_q = 1 \quad (0.12)$$

which gives

$$\frac{\partial I}{\partial p}\Big|_q \left[\frac{\partial\varphi}{\partial q}\Big|_p + \frac{\partial p}{\partial q}\Big|_I \frac{\partial\varphi}{\partial p}\Big|_q \right] = 1 \quad (0.13)$$

or

$$\frac{\partial p}{\partial I}\Big|_q = \left(\frac{\partial I}{\partial p}\Big|_q \right)^{-1} = \frac{\partial\varphi}{\partial q}\Big|_p + \frac{\partial p}{\partial q}\Big|_I \frac{\partial\varphi}{\partial p}\Big|_q = \frac{\partial\varphi}{\partial q}\Big|_I \quad (0.14)$$

and this is what we wanted to prove, see eq. (0.9). In any case, this relation is easier found using a generating function, when

$$\frac{\partial\varphi}{\partial q}\Big|_I = \frac{\partial^2 W(q, I)}{\partial q \partial I} = \frac{\partial p}{\partial I}\Big|_q \quad (0.15)$$

Returning to (0.8) we get

$$\Delta\varphi = \oint \frac{\partial p}{\partial I}\Big|_q dq = \frac{\partial}{\partial I} \oint p dq = 2\pi \frac{\partial I}{\partial I} = 2\pi \quad (0.16)$$