

Computer Simulations in Physics

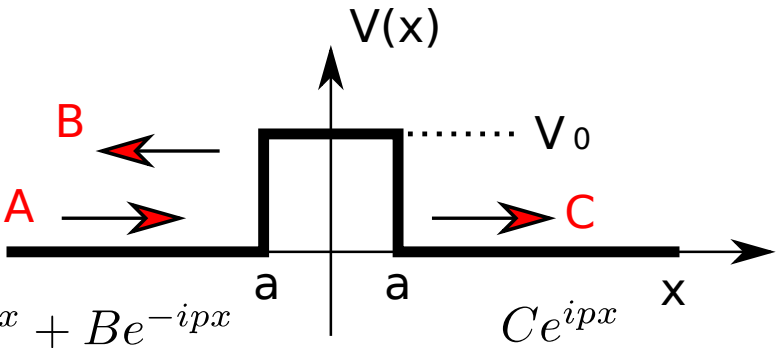
Course for MSc physics students

Solving the Schrödinger equation for 1-dimensional scattering

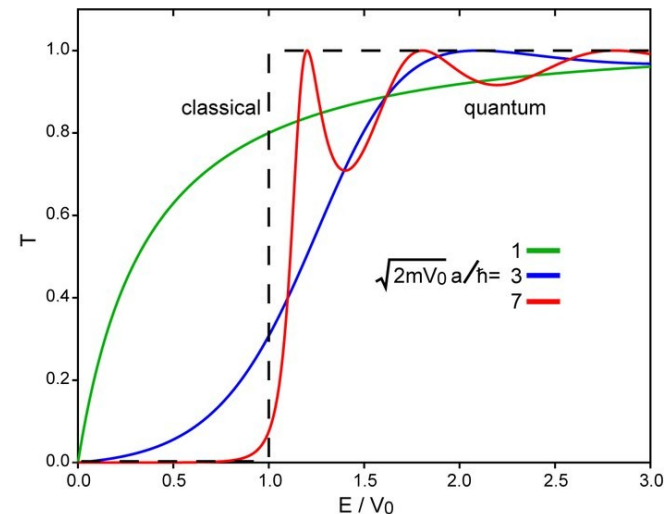
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$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + V(x)\psi(x) = E\psi(x)$$


A diagram of a 1D potential barrier $V(x)$ of height V_0 and width $2a$, centered at $x=0$. The potential is zero for $|x| > a$ and V_0 for $|x| < a$. An incident wave A (red arrow) moves from left to right in the region $x < -a$. A reflected wave B (red arrow) moves from left to right in the region $x < -a$. A transmitted wave C (red arrow) moves from left to right in the region $x > a$. The wave function in the region $x < -a$ is labeled $Ae^{ipx} + Be^{-ipx}$ and in the region $x > a$ is labeled Ce^{ipx} .



Literature

- Numerics: Computational Quantum Physics course at ETH Zurich SS 2008, by P. de Forcrand & M. Troyer
 - lecture notes online
- Quantum Scattering Theory: Any introductory Quantum Mechanics book, e.g., Griffiths

1-dimensional Quantum Mechanics, brief reminder

state of particle:
complex valued wavefunction

$$\Psi(x)$$

position probability density:

$$|\Psi(x)|^2$$

expectation value of position:

$$\langle x \rangle = \int dx |\Psi(x)|^2 x = \int dx \Psi(x)^* x \Psi(x)$$

position operator: $\hat{x} = x \cdot$ $\hat{x}\Psi(x) = x\Psi(x)$

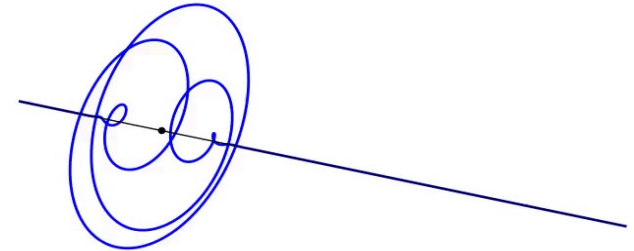
momentum: $\hat{p} = -i\hbar\partial_x$ $\langle p \rangle = \int dr \Psi(x)^* \hat{p}\Psi(x)$

Hamiltonian = operator of total energy: $\hat{H} = \frac{\hat{p}^2}{2m} + V$

time evolution, Schrodinger equation: $i\hbar\partial_t\Psi = \hat{H}\Psi$

$$i\hbar\partial_t\Psi(x) = -\frac{\hbar^2}{2m}\partial_x^2\Psi(x,t) + V(x)\Psi(x,t)$$

Gaussian wave packet, $a = 2$, $k_0 = 4$



1-dimensional quantum mechanics: eigenstates of Hamiltonian, time-independent Schrodinger equation

$$i\hbar\partial_t\Psi(x) = -\frac{\hbar^2}{2m}\partial_x^2\Psi(x,t) + V(x)\Psi(x,t)$$

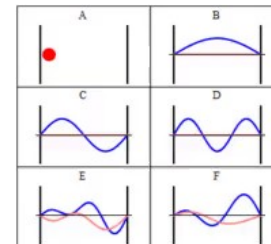
time evolution, Schrodinger equation: $i\hbar\partial_t\Psi = \hat{H}\Psi$

try it on eigenfunction
of the operator H:

$$\hat{H}\psi(x) = E\psi(x) \quad \rightarrow \quad \Psi(x,t) = e^{-iE/\hbar t}\psi(x)$$

This Ψ is a *stationary state*, position distribution $|\Psi(x,t)|^2$ independent of time

Example:
bound states in
a square well

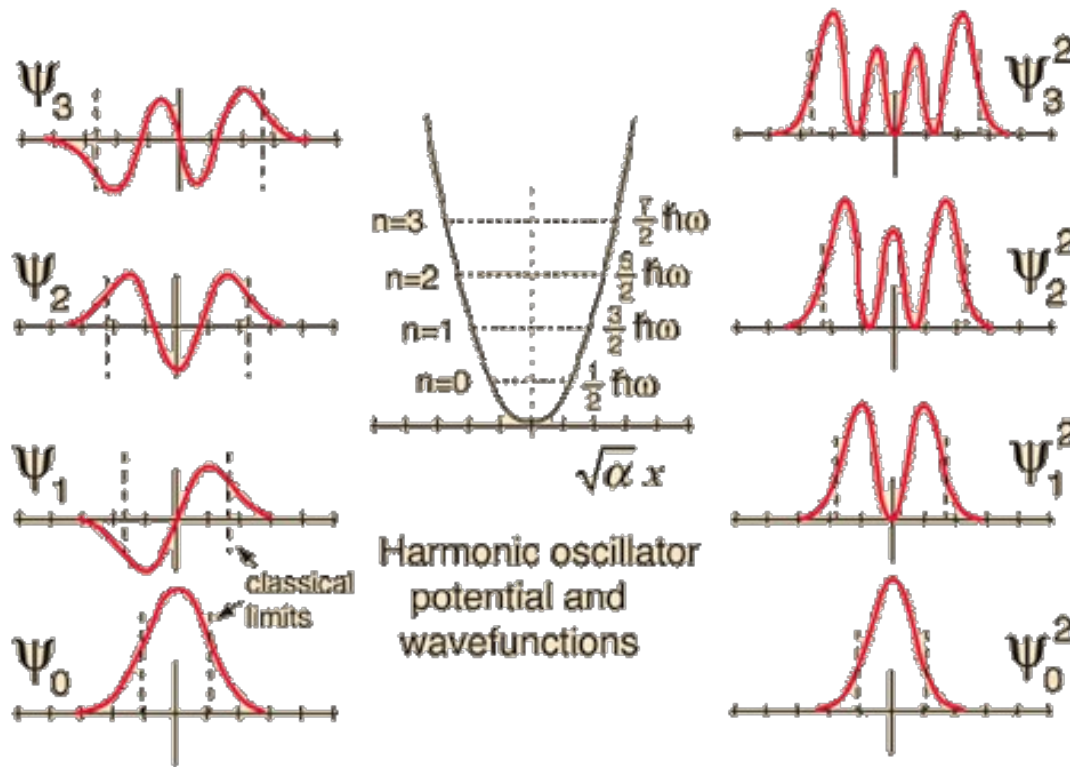


Some trajectories of a particle in a box according to [Newton's laws of classical mechanics](#) (A), and according to the [Schrödinger equation of quantum mechanics](#) (B-F). In (B-F), the horizontal axis is position, and the vertical axis is the real part (blue) and imaginary part (red) of the [wavefunction](#). The states (B,C,D) are [energy eigenstates](#), but (E,F) are not.

Bound states in 1D have real-valued wavefunctions with compact support

$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + \underbrace{\frac{m\omega^2}{2}x^2}_{V(x)}\psi(x) = E\psi(x)$$

$\Psi(x, t) \xrightarrow{x \rightarrow \pm\infty} 0$ particle does not escape



First four harmonic oscillator normalized wavefunctions

$$\Psi_0 = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-y^2/2}$$

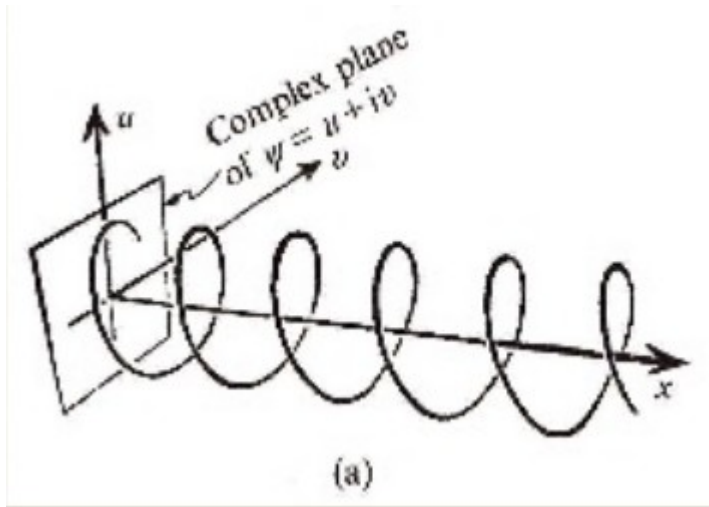
$$\Psi_1 = \left(\frac{\alpha}{\pi}\right)^{1/4} \sqrt{2}y e^{-y^2/2}$$

$$\Psi_2 = \left(\frac{\alpha}{\pi}\right)^{1/4} \frac{1}{\sqrt{2}}(2y^2 - 1)e^{-y^2/2}$$

$$\Psi_3 = \left(\frac{\alpha}{\pi}\right)^{1/4} \frac{1}{\sqrt{3}}(2y^3 - 3y)e^{-y^2/2}$$

$$\alpha = \frac{m\omega}{\hbar} \quad y = \sqrt{\alpha}x$$

Freely propagating particle: plane wave, wavepacket



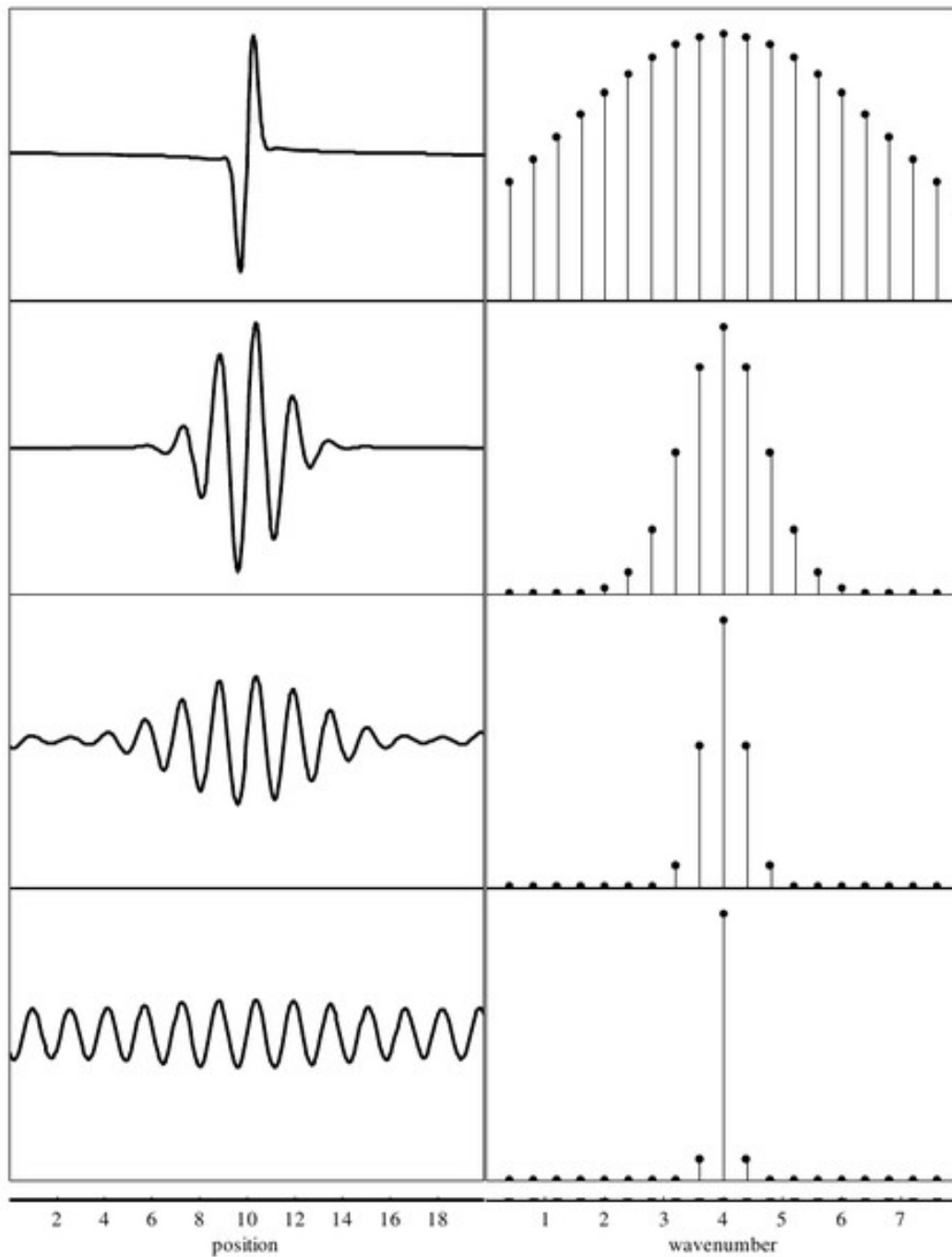
$$\Psi(x, t) = e^{ikx - iE/\hbar t}$$

Wavefunction not normalized for probability but for particle current

State of actual particle: wavepacket

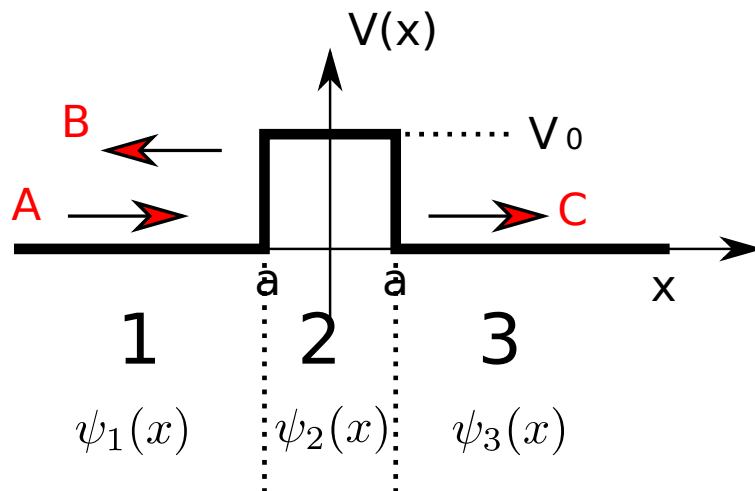


Real space – momentum space: Fourier transform



Scattering states in 1D: incoming, reflected, transmitted plane wave + something in the middle

$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + V(x)\psi(x) = E\psi(x)$$



scattering state:

- a solution of Schrodinger eqn,
= eigenstate of H
- with no incoming wave from right ($D=0$)

→ Away from scattering region (1,3):
superposition of plane waves
with wavenumber k_0

$$k_0 = \sqrt{2mE/\hbar^2}$$

→ In scattering region (2):
depends on potential

$$\psi_1(x) = Ae^{ik_0x} + Be^{-ik_0x} \quad \psi_3(x) = Ce^{ik_0x} + De^{-ik_0x}$$

Solution over all x : fit solutions at a and $-a$:

$$\begin{aligned} \psi_1(-a) &= \psi_2(-a) & \psi_2(a) &= \psi_3(a) \\ \psi'_1(-a) &= \psi'_2(-a) & \psi'_2(a) &= \psi'_3(a) \end{aligned}$$

reflection & transmission amplitudes:

$$r = \frac{B}{A}; \quad t = \frac{C}{A}$$

reflection & transmission probabilities: $R = |r|^2 \quad T = |t|^2$

Example: scattering from Square barrier

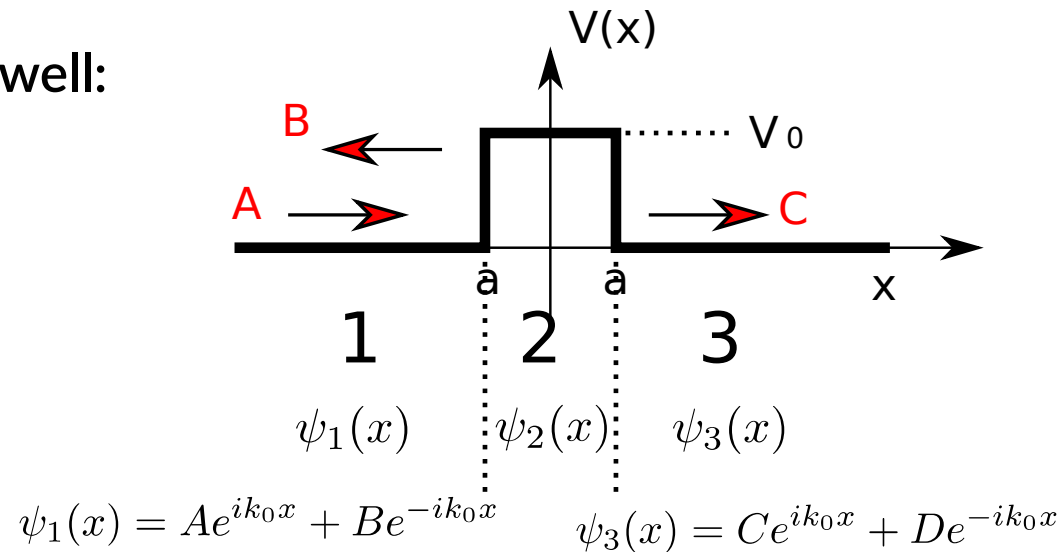
Exactly solvable textbook problem (Griffiths Quantum Mechanics 2.6., or wikipedia, plane waves+fitting)

Solution simple in scattering region as well:

$$\psi_2(x) = Fe^{ik_1x} + Ge^{-ik_1x}$$

$$k_1 = \sqrt{2m(E - V_0)/\hbar^2}$$

$$k_0 = \sqrt{2mE/\hbar^2}$$



Solution over all x: fit solutions at a and -a:

$$\psi_1(-a) = \psi_2(-a)$$

$$\psi_2(a) = \psi_3(a)$$

$$\psi'_1(-a) = \psi'_2(-a)$$

$$\psi'_2(a) = \psi'_3(a)$$

reflection & transmission amplitudes:

$$r = \frac{B}{A}; \quad t = \frac{C}{A}$$

reflection & transmission probabilities: $R = |r|^2 \quad T = |t|^2$

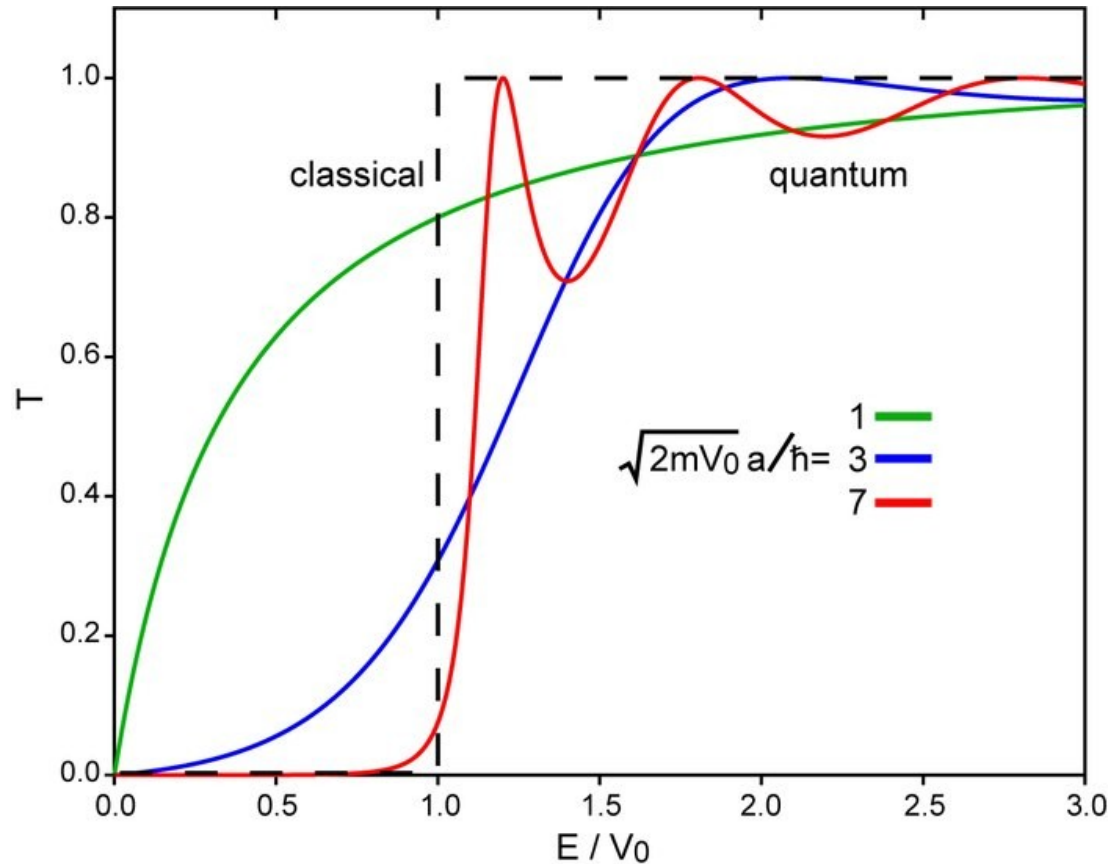
Transmission by tunneling and resonances in square barrier

Tunneling across barrier in classically forbidden regime:

$$T = \frac{1}{1 + \frac{V_0^2}{4E(E-V_0)} \sinh^2 2k_1 a}$$

Transmission across barrier in classically allowed regime:

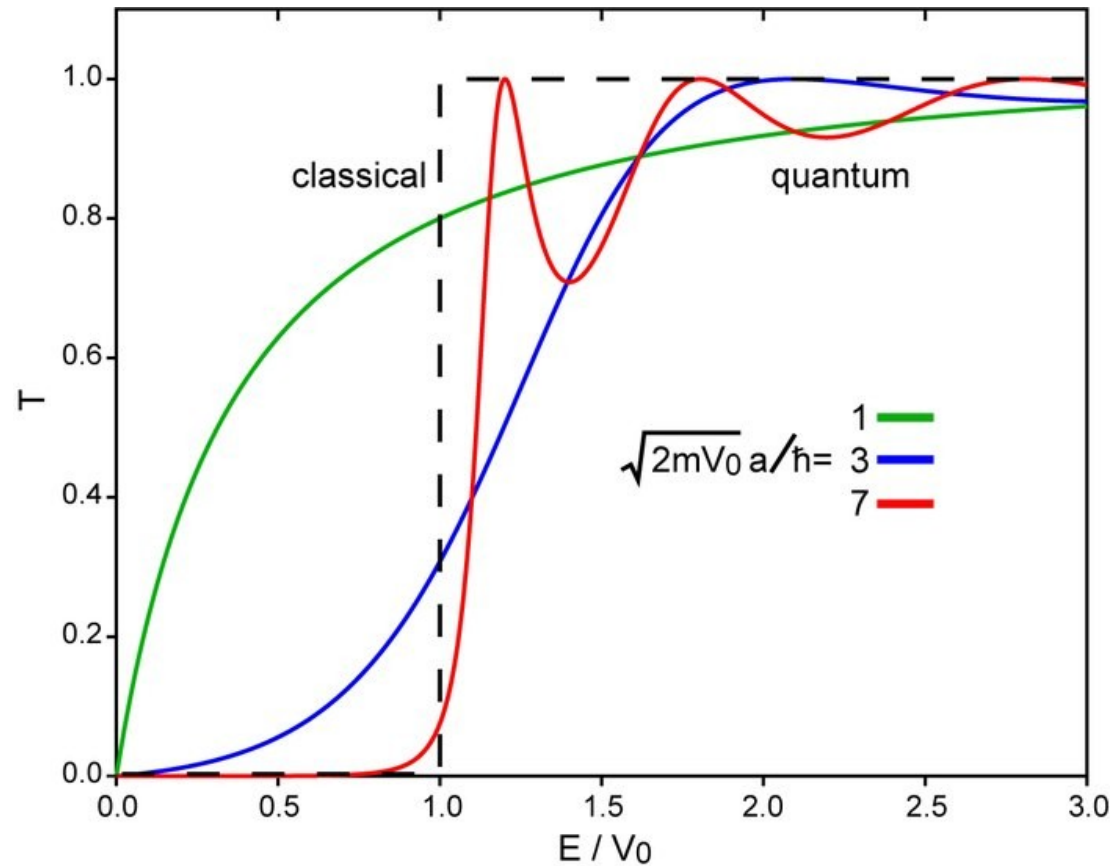
$$T = \frac{1}{1 + \frac{V_0^2}{4E(E-V_0)} \sin^2 2k_1 a}$$



Transmission resonances:
perfect transmission if $a = n \lambda / 2$

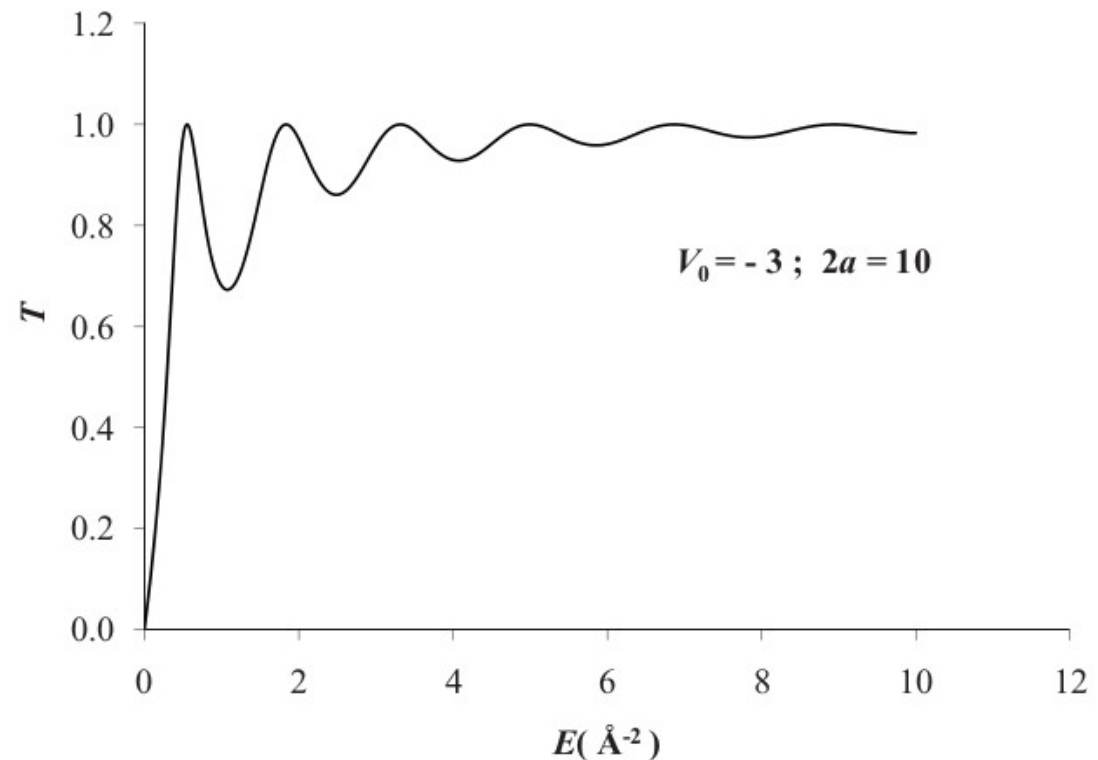
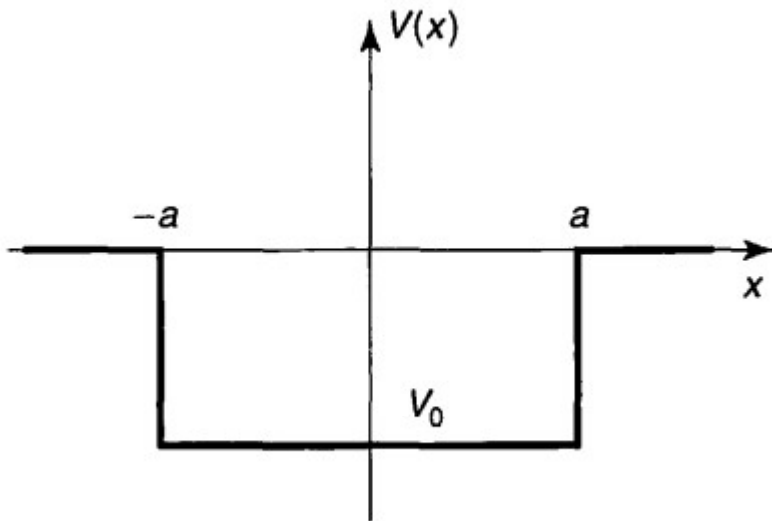
$$E_n = V_0 + \frac{\pi^2 \hbar^2}{8ma^2} n^2$$

What is fundamental in quantum mechanics, what is only for square well?



Formulas also hold for transmission across square well, transmission resonances

As before, with $V_0 < 0$.



Transmission across square well is always classically allowed:

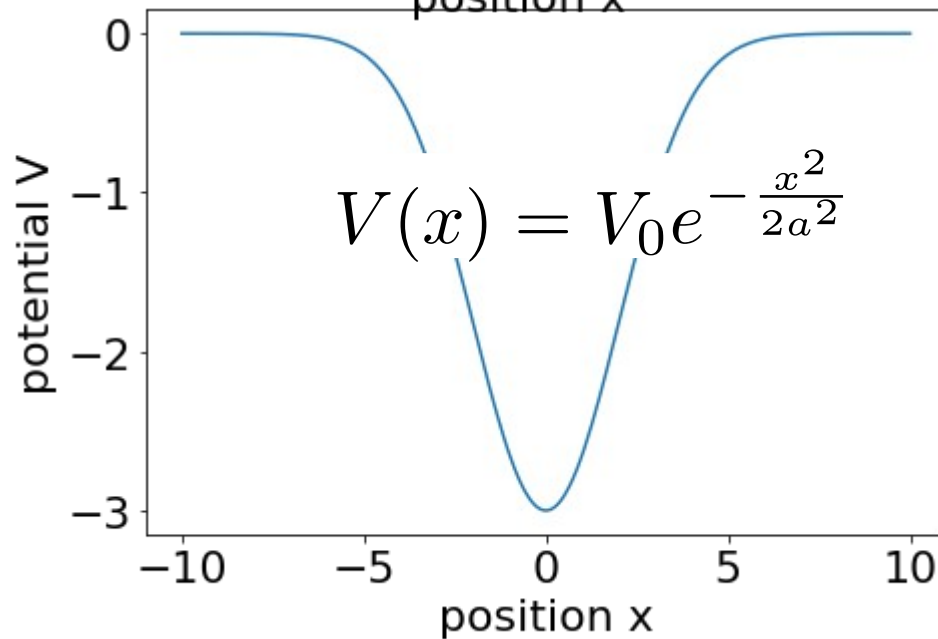
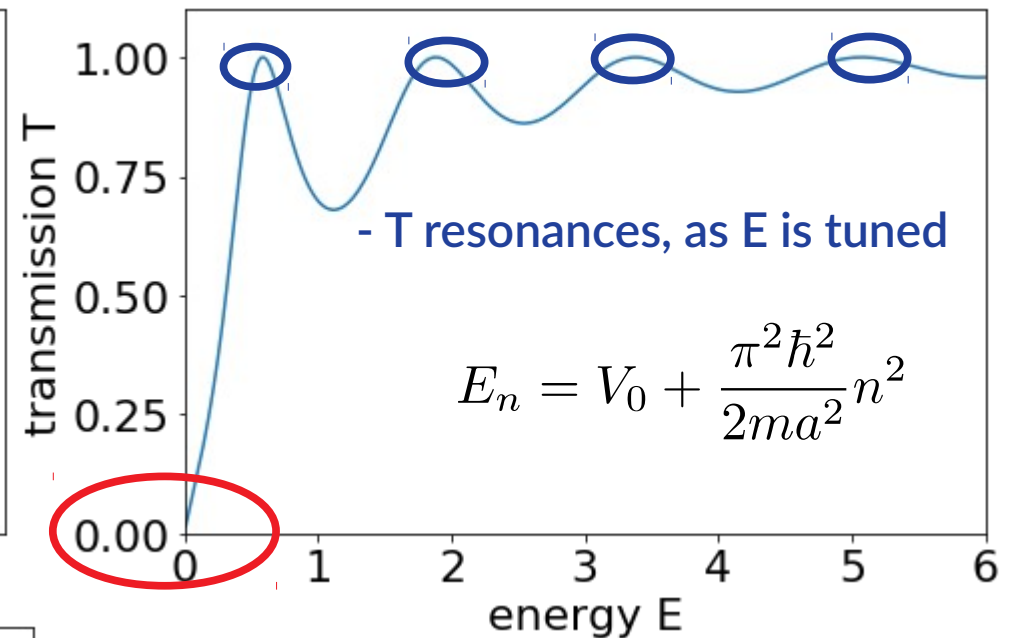
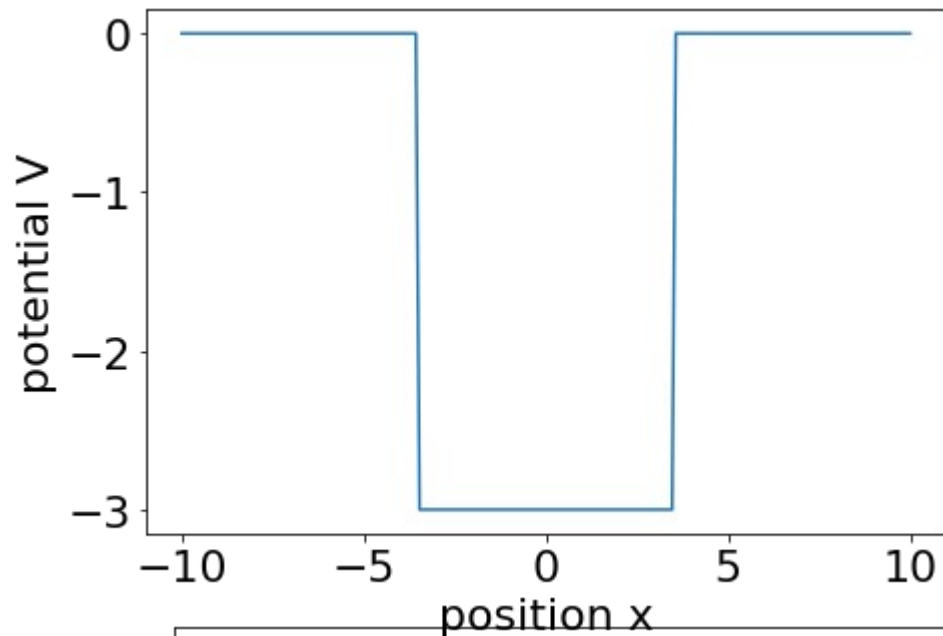
$$T = \frac{1}{1 + \frac{V_0^2}{4E(E - V_0)} \sin^2 2k_1 a}$$

No transmission for $E \rightarrow 0$

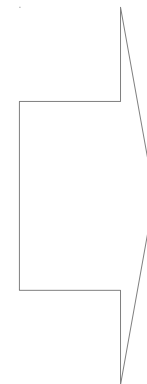
$$E_n = V_0 + \frac{\pi^2 \hbar^2}{8ma^2} n^2$$

What about a transmission resonance at 0 energy?

What happens to transmission resonances in smooth potentials?



- T resonances at E=0, as V0 is tuned?



- T resonances, as E is tuned?

- T resonances at E=0, as V0 is tuned?

Numerov: finite-difference method to solve Schrodinger equation (like Runge-Kutta)

$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + V(x)\psi(x) = E\psi(x)$$

discretize position: $x_n = n\Delta x$ $\psi_n = \psi(x_n)$

Taylor expand ψ :

$$\psi_{n\pm 1} = \psi_n \pm \Delta x \psi'_n + \frac{\Delta x^2}{2} \psi''_n \pm \frac{\Delta x^3}{6} \psi_n^{(3)} + \frac{\Delta x^4}{24} \psi_n^{(4)} \pm \frac{\Delta x^5}{120} \psi_n^{(5)} + O(\Delta x^6)$$

Use a trick to get rid of all odd order derivatives:

$$\psi_{n+1} + \psi_{n-1} = 2\psi_n + (\Delta x)^2 \psi''_n + \frac{(\Delta x)^4}{12} \psi_n^{(4)}.$$

Approximate 4th derivative as a finite difference:

$$\psi_n^{(4)} = \frac{\psi''_{n+1} + \psi''_{n-1} - 2\psi''_n}{\Delta x^2}$$

Numerov method, summarized

$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + V(x)\psi(x) = E\psi(x)$$

dimensionless variables:

$$\psi''(x) + k(x)\psi(x) = 0$$

$$\left(1 + \frac{(\Delta x)^2}{12}k_{n+1}\right)\psi_{n+1} = 2\left(1 - \frac{5(\Delta x)^2}{12}k_n\right)\psi_n - \left(1 + \frac{(\Delta x)^2}{12}k_{n-1}\right)\psi_{n-1} + O(\Delta x^6)$$

locally accurate to 5th order

To calculate with Numerov method, need initial conditions: 2 neighboring values, to iterate

If potential is finite range, $V(x) = 0$ for $|x| > a$
→ use plane wave/decaying form

$$\psi(-a) = 1$$

(bound state at -E):

$$\psi(-a - \Delta x) = \exp(\Delta x \sqrt{2mE/\hbar})$$

scattering state at +E:

$$\psi(-a - \Delta x) = \exp(\pm i \Delta x \sqrt{2mE/\hbar})$$

Scattering problem by Numerov algorithm

$$\psi_L(x) = A \exp(-iqx) + B \exp(iqx)$$

$$\psi_R(x) = C \exp(-iqx)$$

- Set $C = 1$ and use the two points a and $a + \Delta x$ as starting points for a Numerov integration.
- Integrate the Schrödinger equation numerically – backwards in space, from a to 0 – using the Numerov algorithm.
- Match the numerical solution of the Schrödinger equation for $x < 0$ to the free propagation ansatz (3.11) to determine A and B .

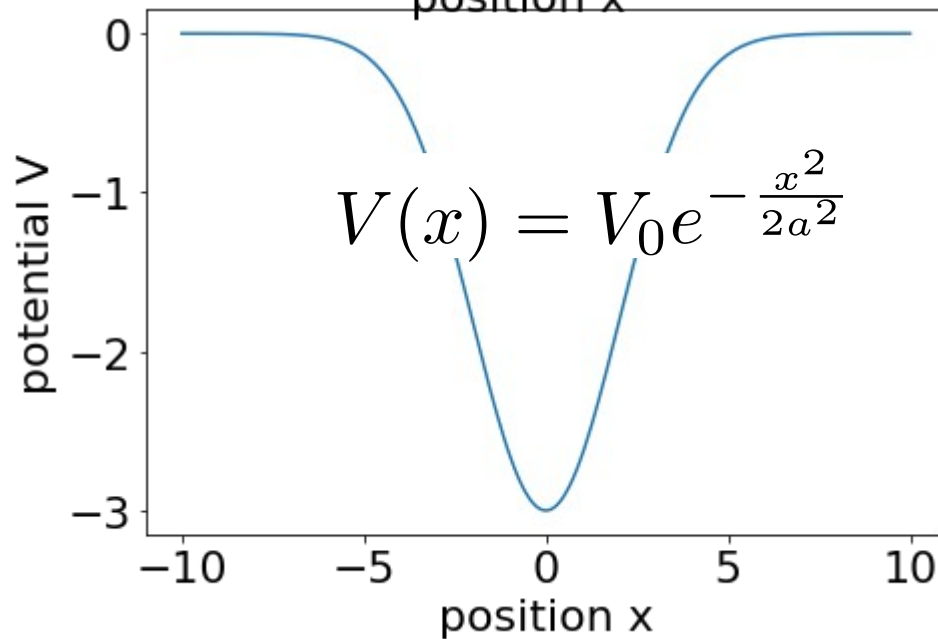
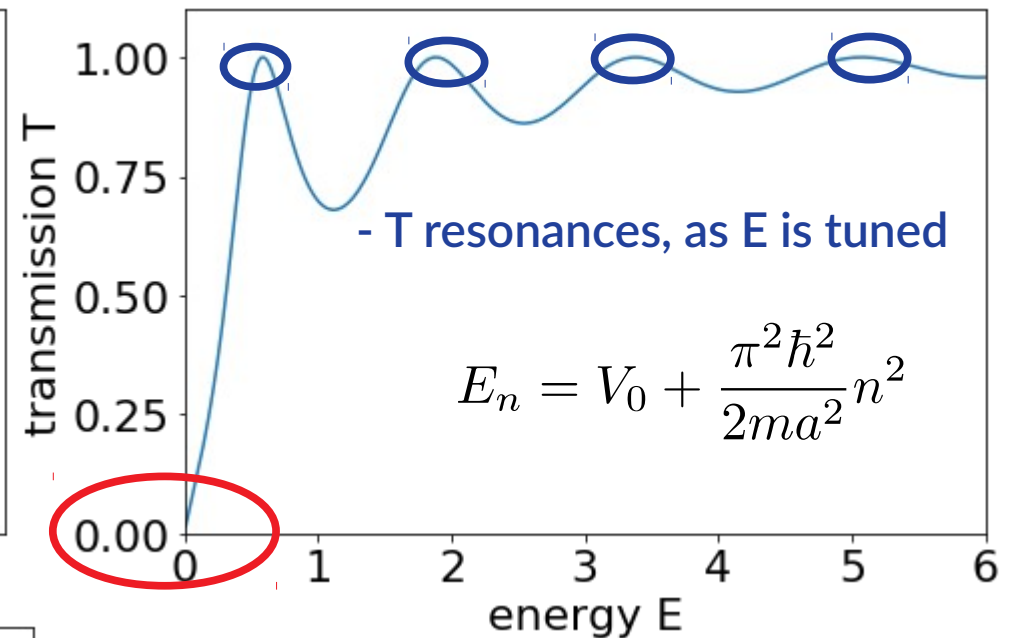
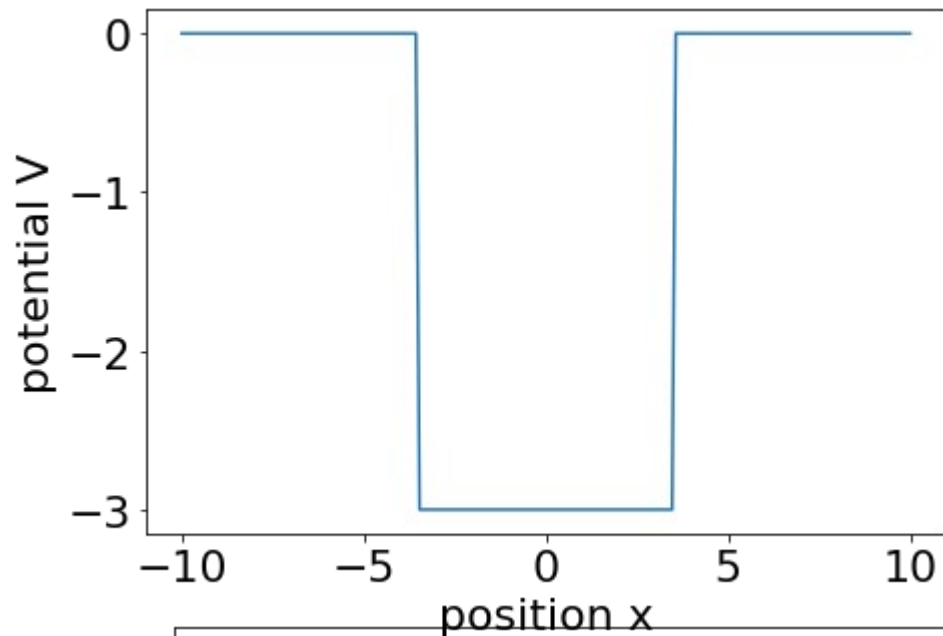
$$R = |B|^2/|A|^2$$

$$T = 1/|A|^2$$

Exercises for today

- Calculate scattering from a square potential barrier by integration of Schrodinger equation using Numerov
 - Plot transmission as a function of energy, for a barrier height 2 eV, size 1nm
 - Compare with analytical curves
- Calculate scattering from a Gaussian potential barrier by the same method
 - Plot transmission as before, with a barrier height 2eV, size 1nm
 - What happened to the transmission resonances?

Homework: what happens to transmission resonances in smooth potentials?



- T resonances at E=0, as V0 is tuned?

- T resonances, as E is tuned?

- T resonances at E=0, as V0 is tuned?

Homework exercises

- Calculate scattering from a square potential well by integration of Schrodinger equation using Numerov
 - Plot transmission as a function of energy, for a well depth 2 eV, size $2a=1\text{nm}$.
 - Also plot analytical curves
 - Plot transmission as a function of well depth (4 eV \rightarrow 0 eV), at energy 0.1 eV, well size $a=0.5\text{ nm}$
- Change the shape of the potential well to Gaussian. What happens to the resonances in the two cases above?
 - Plot transmission as a function of energy, for a well depth 2 eV, size $a=0.5\text{nm}$
 - Plot transmission as a function of well depth (4 eV \rightarrow 0 eV), at energy 0.1 eV, well size $a=0.5\text{nm}$