



# **Laser Physics 11.**

## **Laser basics, fundamental knowledge about laser operation**

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Date	Topic	Lecture Seminar	Date	Topic	Lecture Seminar
24-Oct	Introduction Light-matter interaction	L	26-Oct	Interaction of light with matter – line broadening	L-S
31-Oct (5-Nov)	Coherent optical amplifier (2)–saturation and noise	L	2_Nov	Coherent optical amplifier (1)– gain, bandwidth, phase	L-S
7-Nov	Fabry-Perot resonators, longitudinal modes	L	9-Nov	Passive optical resonators, stability, modes	S
14-Nov	Properties of Gaussian beams – physical optics	S	16-Nov	break	
21-Nov	<b>Control test</b>	S	23-Nov	Gain conditions, phase condition spectral bandwidth	L
28-Nov	Pulsed mode operation, Q- switching, mode-locking	L	30-Nov	Properties of laser beams – coherence, monochromaticity, brightness	L
5-Dec	Properties of laser beams - spatial coherence, divergence	L-S	7-Dec	Semiconductor lasers and amplifiers, homo and heterojunction, properties	L-S



# LASER - acronym

Light **A**mplification by **S**timulated **E**mission of **R**adiation

Stimulated emission – concept: by A. Einstein, **1917**

A. E. deduced the energy density (energy per unit volume and frequency interval) for the emission of electromagnetic radiation of two-level atoms in thermal equilibrium as a function of the temperature (Planck's law of black-body radiation):

$$\rho(\nu) = \frac{8\pi}{c^3} \nu^2 h\nu \frac{1}{e^{h\nu/kT} - 1}$$

First laser - ruby laser in **1960**

**What are the special properties of a laser?**



## Special properties:

Directed beam with small divergence

(1 mrad = 3.4 angular minute),

except for diode lasers ( $10\text{-}30^\circ$  divergence)

Monochromatic or narrow spectral bandwidth light: 0.1 Hz

for short time, 50-100 Hz long-time operation is feasible

Coherent beam: spatial and temporal coherence

Ultra short laser pulses:  $10^{-15}$  s (fs)



## Special properties:

Power range:  $10^{-9} - 10^{20}$  W

Wavelength range:  $(10^{11} - 10^{17}$  Hz)

Far-infrared: 10 – 1000  $\mu\text{m}$

Mid-infrared: 1 – 10  $\mu\text{m}$

Near-infrared: 0.7 – 1  $\mu\text{m}$

**Visible:** 0.4 – 0.7  $\mu\text{m}$  or 400 – 700 nm ( **$10^{14} - 10^{15}$  Hz**)

UV: 0.2 – 0.4  $\mu\text{m}$  or 200 – 400 nm

Vacuum UV: 0.1 – 0.2  $\mu\text{m}$  or 100 – 200 nm

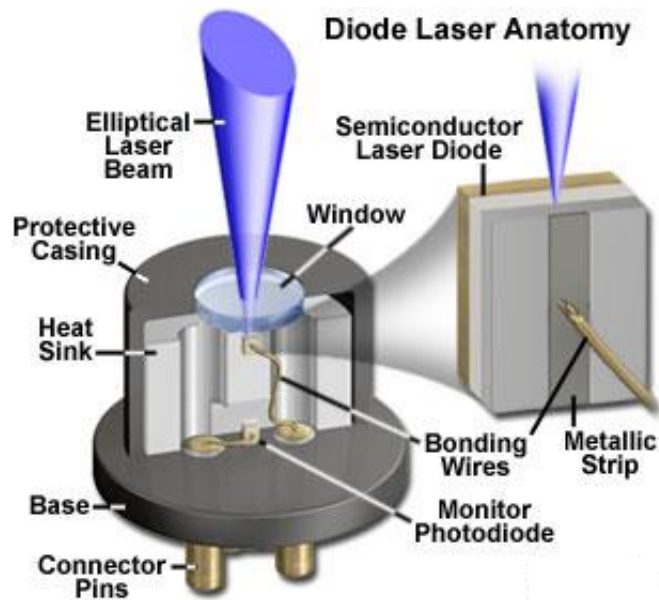
E(Extreme)UV: 10 – 100 nm

Soft x-ray: 1 – 20 (30) nm (overlap with EUV)



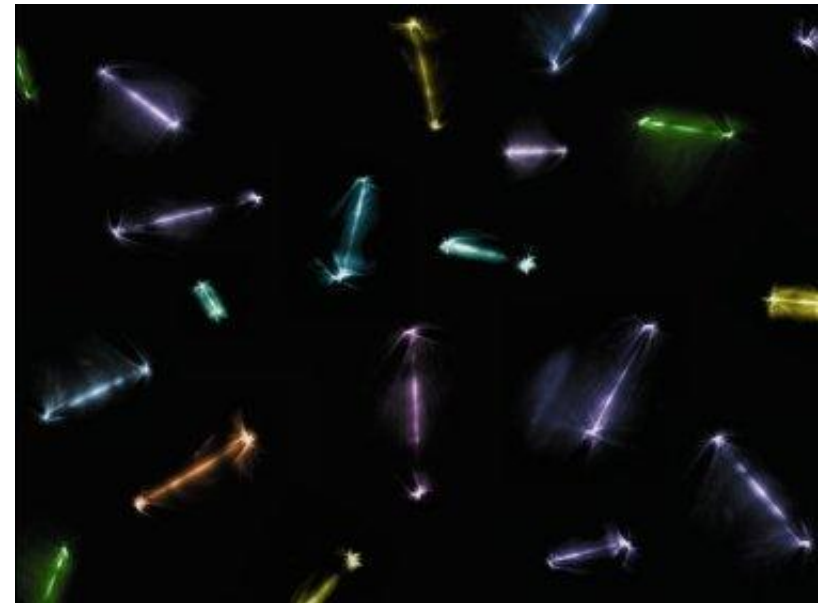
# Special properties:

Dimension - small



Laser diode

$0.4 \mu\text{m} \times 2 \mu\text{m} \times 400 \mu\text{m}$



Multicolor nanowire lasers

$\text{Ø} 200 - 400 \text{ nm}$ , length 20 – 60  $\mu\text{m}$

Source: Optics & Laser Europe, August, 2008



# Special properties:

Dimension - large



NIF, Lawrence Livermore National Laboratory, laser system with 192 beams: 500 TW, 7456 m laser



# Milestones in the laser history

Date	Event	Person/institute
<b>1917</b>	<b>Concept of the stimulated emission</b>	<b>Albert Einstein</b>
1951	MASER (Nobel prize in 1964)	Townes, Prohorov, Basov
<b>1960</b>	laser patent 2.929.922	Shawlow and Townes
	<b>first laser – ruby, solid state laser</b>	Th. Maiman, Hughes Labs.
1961	first gas laser He-Ne (infrared, red in 1962)	Bell Laboratories
1962	Pulsed semiconductor laser (GaAs) at liquid N <sub>2</sub> temperature	General Electric
	Ophthalmological applications (rubin and He-Ne)	
1963	first CO <sub>2</sub> laser	K. N. Patel, Bell Labs.
1964	first Nd-YAG laser	Bell Laboratories
	first Ar-ion laser	Hughes Laboratories
	first surgical experiments on animals	
	distance measurement with a laser interferometer	
1965	first chemical laser	California Berkeley Univ.
	first surgical CO <sub>2</sub> laser	
	<b>proposal - fiber optic communications (Nobel prize 2009)</b>	
1966	first dye laser	IBM Laboratories
1968	Ar-ion laser application in ophthalmology and urology	
1969	First three lasers for material processing in the car industry	General Motors
	rubin laser distance meter (measurement of the Earth – Moon distance using the mirror installed by the Apollo 11 astronauts on the Moon)	US army
	concept of integrated optics	Stewart Miller



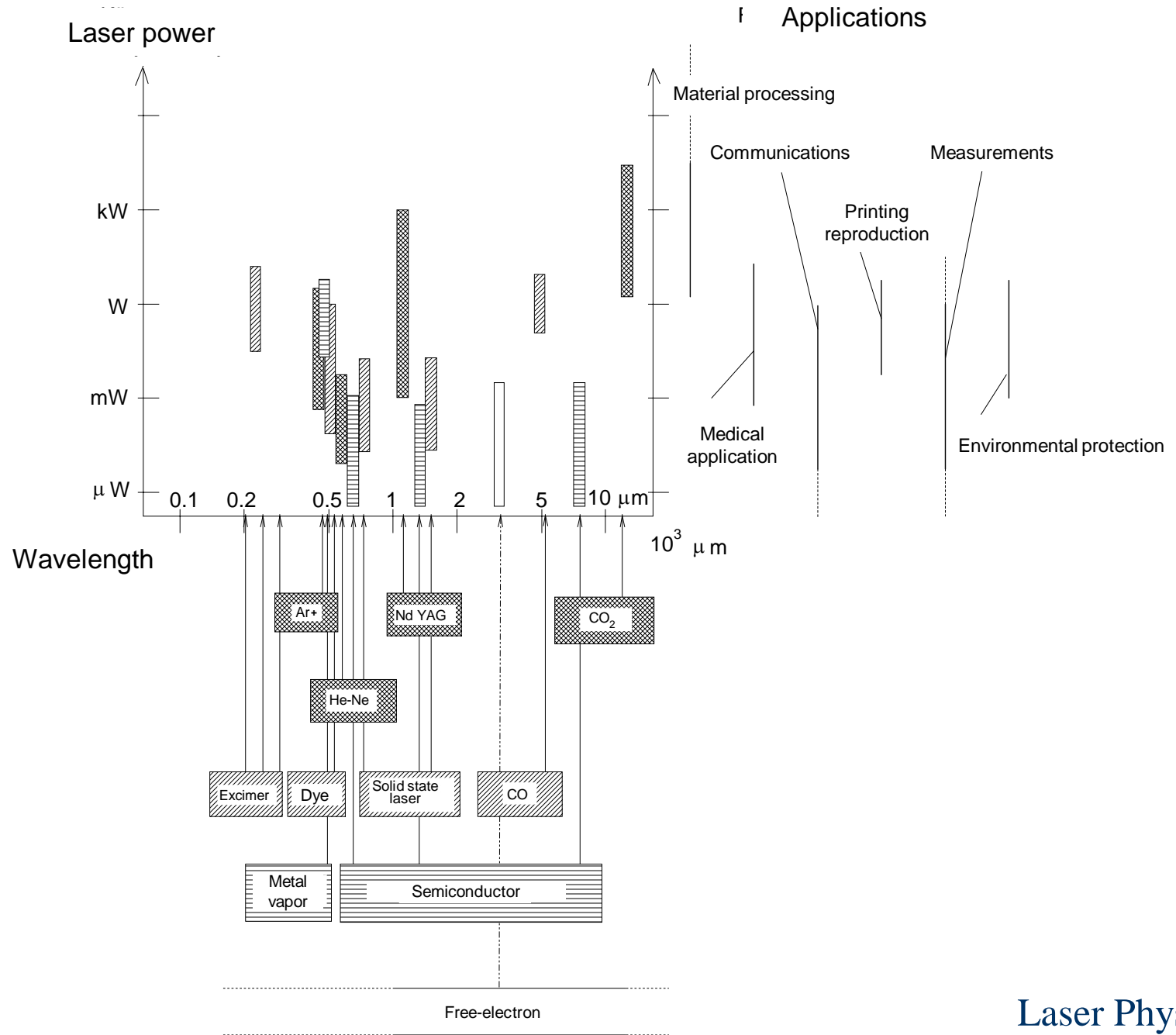


# Milestones in the laser history

1970	first continuous wave (cw) GaAlAs semiconductor laser ( $T_{\text{room}}$ ) 20 dB/km	Corning
1971	first N <sub>2</sub> laser Nobel-prize for the development of holography (principle was published in 1948)	Gábor, Dénes
1972	laser target designator: military use in Vietnam fiber with 4 dB/km loss	US army Corning
1973	Argon and Nd-YAG laser coupled to fiber optic endoscope	Nath, Gorish, Kiefbacher
1975	commercial cw semiconductor laser (laser diode - LD) first medical laser symposium	
1984	first laboratory x-ray laser	Lawrence Livermore Lab.
1987	LD-pumped Nd:YAG laser Er-doped fiber amplifier dye laser with 6 fs pulse length	
1992	20 Gbit/s transmission speed for 100 km	
1995-	Green, blue light emitting diode (LED) laser pulse with petawatt ( $10^{12}$ W) power 23 W fiber laser 5 fs Ti-Sapphire laser pulse	
1999	LD of 400 nm	Nichia



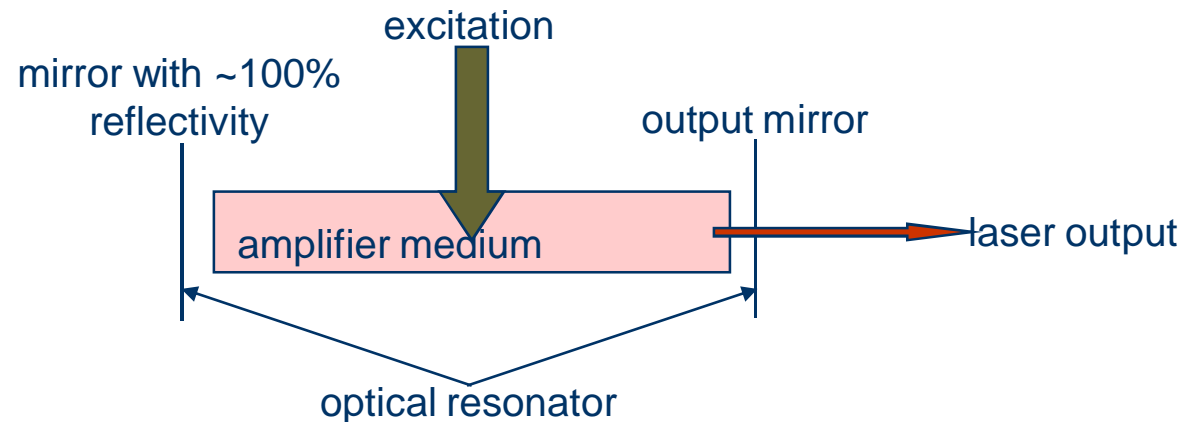
# Types of lasers





# What are needed for a laser?

1. Amplifier medium (gas, liquid, solid state) in which, due to excitation (by light, electric discharge or current, chemical reaction, ...) light amplification by stimulated emission is possible
2. Feedback system – optical resonator





# Selection of the laser material

Particle with known energy structure (atom, molecule, doping ion in a solid)

Selection of levels  $E_1$  and  $E_2$ , in between light emission and absorption are possible

$$E_2 - E_1 = h\nu_0,$$

$\nu_0$  is the resonance frequency.  $E_2$  is an excited state, however both  $E_1$  and  $E_2$  can be excited states. In practice, photon emission or absorption is possible in a range of frequencies around  $\nu \approx \nu_0$ .

3 kinds of interaction:      spontaneous emission,  
   absorption  
   stimulated emission



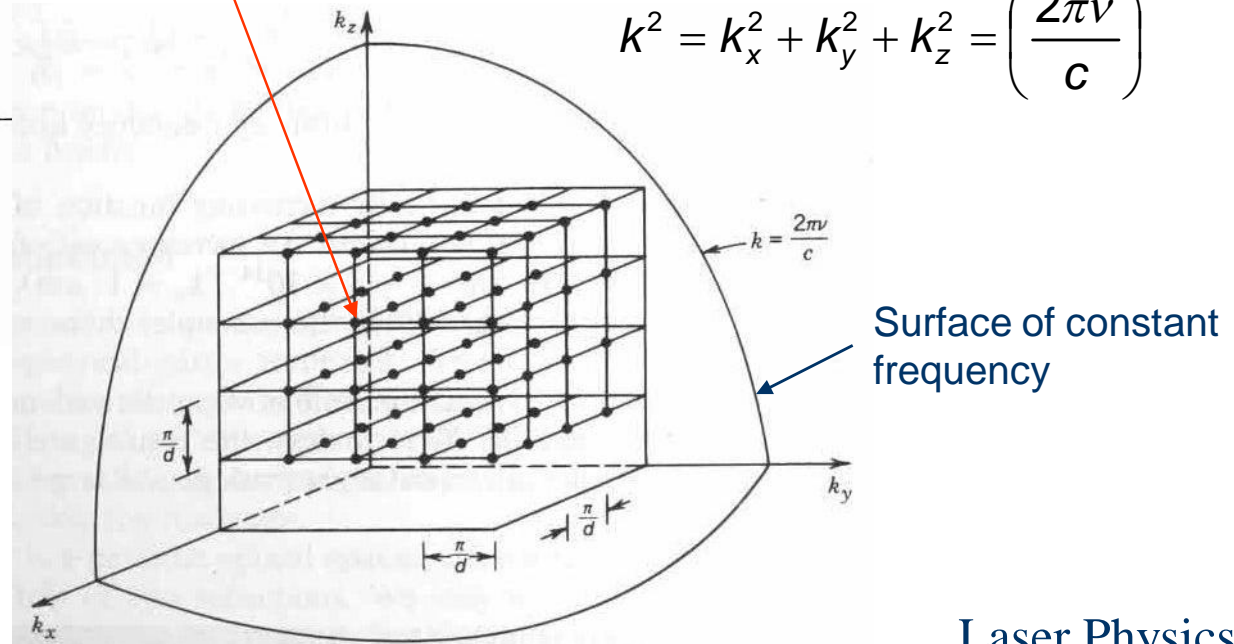
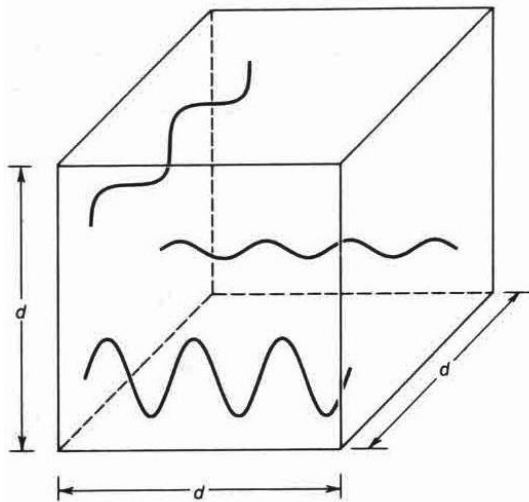
# Laser material in a 3D optical resonator

We place the selected laser material into an optical resonator of volume  $V$ , where  $V = d^3$  ( $d$ : edge of cube with reflecting surfaces). Standing wave solutions:  $\mathbf{k} = (k_x, k_y, k_z)$  discrete values.

$$k_x = \frac{q_x \pi}{d}, \quad k_y = \frac{q_y \pi}{d}, \quad k_z = \frac{q_z \pi}{d}, \quad q_x, q_y, q_z = 1, 2, \dots$$

Possible frequencies :

$$k^2 = k_x^2 + k_y^2 + k_z^2 = \left( \frac{2\pi\nu}{c} \right)^2$$





# Laser material in a 3D optical resonator

Number of modes in the  $0 - \nu$  frequency range; density of modes?

The number of modes with continuum approximation (volume of 1/8 sphere / volume of the cell):

$$2 \left( \frac{1}{8} \right) \frac{4\pi k^3 / 3}{(\pi/d)^3} = \left( \frac{k^3}{3\pi^2} \right) d^3, \quad k = \frac{2\pi\nu}{c}$$

due to polarization

$$\frac{8\pi^3\nu^3}{3c^3\pi^2} d^3 = \frac{8\pi\nu^3}{3c^3} d^3$$

Density of modes (the number of modes in unit volume and frequency interval):

$$M(\nu) = \frac{8\pi\nu^2}{c^3}$$



# Light-matter interactions

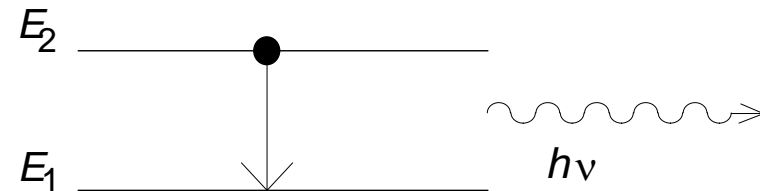
Laser material in an optical resonator of volume  $V$ .  
 $\nu \approx \nu_0$  is the selected mode of the resonator.

We are interested in the possible interactions of the particle with the resonator mode of frequency  $\nu$ .

## Spontaneous emission

The energy of the particle is initially  $E_2$  (upper level). It may decay spontaneously to the lower energy level.

The probability density (per second) or rate -  $\rho_{sp}$  - is a function of the frequency  $\nu$ .



$$\rho_{sp} = \frac{c}{V} \sigma(\nu) \quad [\text{s}^{-1}],$$

$\sigma(\nu)$ , the transition cross section (with surface dimension) peaks around  $\nu_0$ ,  $c$  is the velocity of light in the medium.



# Light-matter interactions

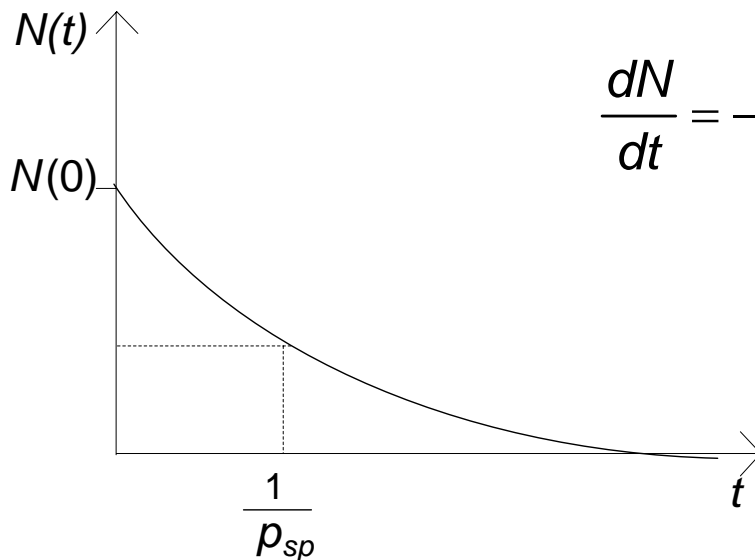
## Spontaneous emission (cont.)

$\rho_{sp}\Delta t$  is the probability, that an emission of photon of frequency  $\nu$  takes place in the time interval of  $[t, t + \Delta t]$ .  $\rho_{sp}\Delta t < 1$ .

If  $N$  is the density of particles in the upper level  $E_2$ , then in unit volume  $\Delta N$  atoms will undergo the transition within  $\Delta t$ :

$$\Delta N = (\rho_{sp}\Delta t)N$$

The rate of change of atoms per unit volume on level  $E_2$ :



$$\frac{dN}{dt} = -\rho_{sp}N, \quad N(t) = N(0)e^{-\rho_{sp}t}.$$

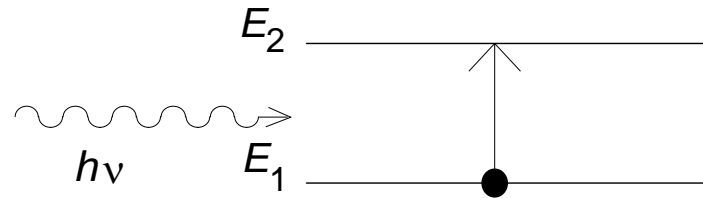




# Light-matter interactions

## Absorption

If the particle is initially on the lower level  $E_1$ , a photon may be absorbed raising the particle to the upper level,  $E_2$ . Absorption is a transition **induced** by the photon.



Effect: the number of photons in frequency mode  $\nu$  decreases in volume  $V$ !

Probability density for the absorption of one photon from the frequency mode  $\nu$  when one photon is in the mode:      when  $n$  photons are in the mode:

$$p_{ab} = \frac{c}{V} \sigma(\nu), \quad P_{ab} = n \frac{c}{V} \sigma(\nu).$$

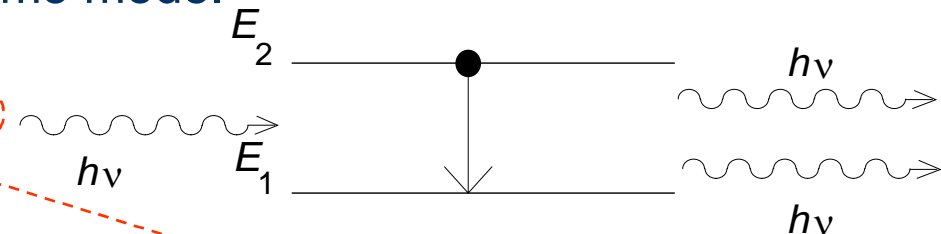


# Light-matter interactions

## Stimulated emission

A particle in level  $E_2$  is stimulated by a photon of frequency  $\nu$  to emit another photon into the same mode.

Effect: the number of photons of frequency  $\nu$  increases with one in  $V$ !



Probability density for one or  $n$  photons:

$$p_{ie} = \frac{c}{V} \sigma(\nu), \quad P_{ie} = n \frac{c}{V} \sigma(\nu).$$
$$P_{ie} = P_{ab} = W_i$$

The clone photon has the same energy, direction and preserves polarization and phase of the mode!!

Spontaneous and stimulated emission together:

$$p_{sp} + P_{ie} = (n + 1) \frac{c}{V} \sigma(\nu).$$



# Light-matter interactions

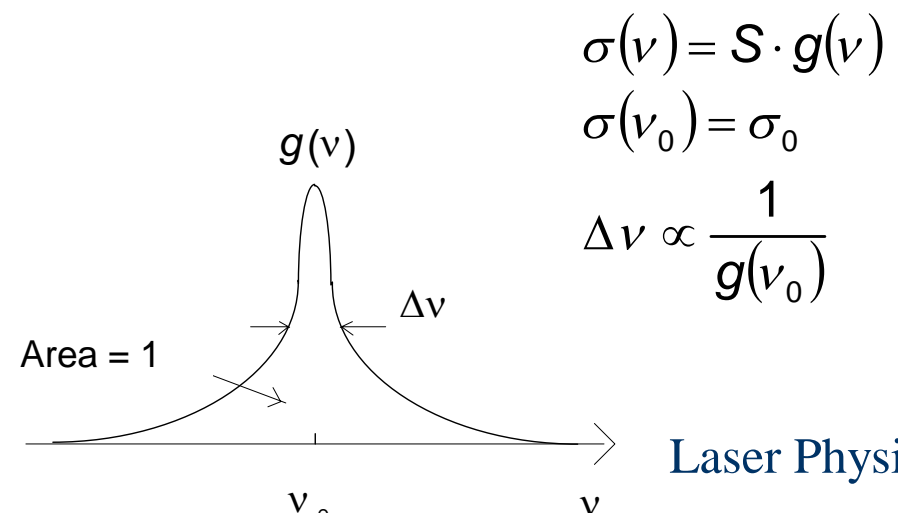
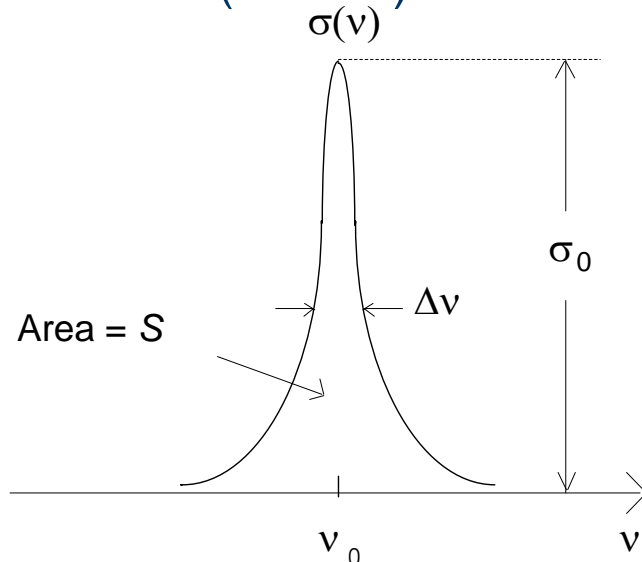
## The lineshape function

The transition cross section  $\sigma(\nu)$  specifies the character of the interactions. Its area represents the strength of the interaction, its shape characterizes the relative frequency dependence of the interaction:

$$S = \int_0^{\infty} \sigma(\nu) d\nu \quad [\text{cm}^2\text{s}^{-1}]$$

$$g(\nu) = \frac{\sigma(\nu)}{S}$$

$g(\nu)$  is the normalized lineshape function.  $\Delta\nu$  is the full width of  $g(\nu)$  at half maximum (FWHM).



$$\sigma(\nu) = S \cdot g(\nu)$$

$$\sigma(\nu_0) = \sigma_0$$

$$\Delta\nu \propto \frac{1}{g(\nu_0)}$$



# Light-matter interactions

## Total spontaneous emission into all modes

Spontaneous emission is possible in  $V$  into all modes with frequency within the spectral width of  $\sigma(\nu)$ :

$$P_{sp} = \int_0^{\infty} \left[ \frac{c}{V} \sigma(\nu) \right] [M(\nu) \cdot V] d\nu = c \int_0^{\infty} \sigma(\nu) M(\nu) d\nu.$$

$\rho_{sp}$       density of modes

$\sigma(\nu)$  is narrow in comparison with  $M(\nu)$ :  $P_{sp} = M(\nu_o) c S = \frac{8 \pi \nu_o^2}{c^2} S = \frac{8 \pi}{\lambda^2} S. \quad [s^{-1}]$

We define a time constant  $t_{sp}$ , known as spontaneous lifetime by:

$$\boxed{P_{sp} = \frac{1}{t_{sp}}} \text{ and } \boxed{S = \frac{\lambda^2}{8 \pi t_{sp}}}.$$

typical value of  $t_{sp}$  is  $10^{-8}$  s, but it can range from  $10^{-13}$  to 100 s.

$$\sigma(\nu) = S g(\nu) = \frac{\lambda^2}{8 \pi t_{sp}} g(\nu)$$

$\lambda$  is the wavelength in the material.



# Light-matter interactions

## Transitions in the presence of broadband light

The spectral energy density in  $V$  is  $\rho(\nu)$  (energy per unit bandwidth per unit volume) is broad as compared to the width of the particle's lineshape function.

The number of photons in the  $[\nu, \nu + d\nu]$  frequency range  $\frac{\rho(\nu)Vd\nu}{h\nu}$ .

$$W_i = \int_0^{\infty} \frac{\rho(\nu)V}{h\nu} \frac{c}{V} \sigma(\nu)d\nu \approx \frac{\rho(\nu_0)}{h\nu_0} c \underbrace{\int_0^{\infty} \sigma(\nu)d\nu}_S = \frac{\rho(\nu_0)}{h\nu_0} cS,$$

We assume that  $\rho(\nu)$  varies slowly in comparison with the sharply peaked  $\sigma(\nu)$  function.

$$W_i = \frac{\lambda^3}{8\pi h t_{sp}} \rho(\nu_0)$$



# Light-matter interactions

## Einstein's $A$ and $B$ coefficients

Based on the analysis of the exchange of energy between atoms and radiation field in thermal equilibrium, Einstein postulated expressions for the probability densities for spontaneous-emission and stimulated transitions:

$$P_{sp} = A$$
$$W_i = B\rho(\nu)$$

$A$  and  $B$  are known as Einstein's  $A$  and  $B$  coefficients. By comparing the above with our expressions:

$$A = \frac{1}{t_{sp}}, \quad B = \frac{\lambda^3}{8\pi h t_{sp}}$$
$$\frac{A}{B} = \frac{8\pi h}{\lambda^3}.$$

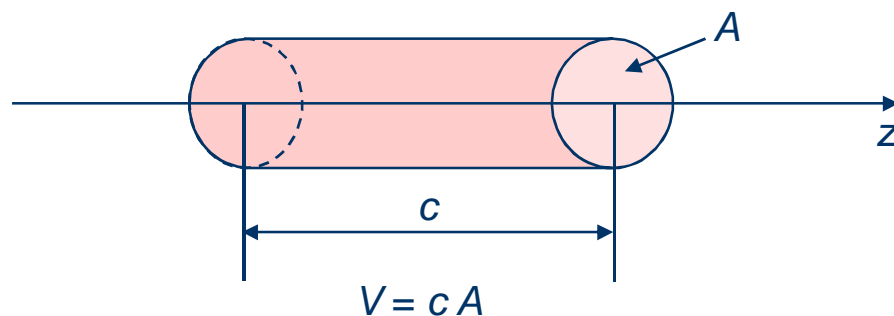


# Light-matter interactions

## Transitions induced by monochromatic light

Interaction of a monochromatic photon beam of frequency  $\nu$  with a particle of resonance frequency  $\nu_0$  (laser type interaction). The photon beam has an intensity  $I$  and travels into the  $z$  direction. The photon-flux density (photons /  $\text{cm}^2 \cdot \text{s}$ ) is:

$$\Phi = \frac{I}{h\nu}. \quad W_i = P_{ab} = P_{ie} \left( = n \frac{c}{V} \sigma(\nu) \right) = ?, \quad n = ?$$



$\Phi A$  photons go through the area  $A$  in unit time. As the length of the cylinder is equal to  $c$  (velocity of light), all photons will leave the cylinder through its base:

$$n = \Phi A = \Phi \frac{V}{c}, \quad \boxed{W_i = \Phi \sigma(\nu)}$$