Introduction to high energy resolution X-ray spectroscopies

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1. Introduction

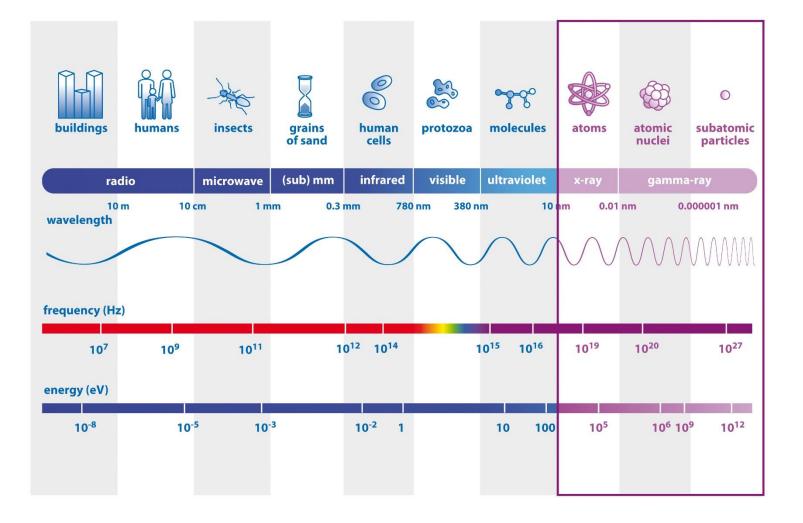
Literature:

 Vankó György: A szinkrotronsugárzás újszerű kémiai alkalmazásai: nagyfelbontású röntgenspektroszkópia - A kémia újabb eredményei 100. (Budapest, 2008)

By X-ray spectroscopies here I mean...

- Typically hard X-rays: 4-20 keV (but: soft X-rays below 1kV and tender regime between soft and hard)
- Absorption and emission of X-ray photons
- High energy resolution: better than the natural broadening of the signal
- Can be used with high spatial resolution
- Can be used with high time resolution
- Can be used to probe magnetic polarization

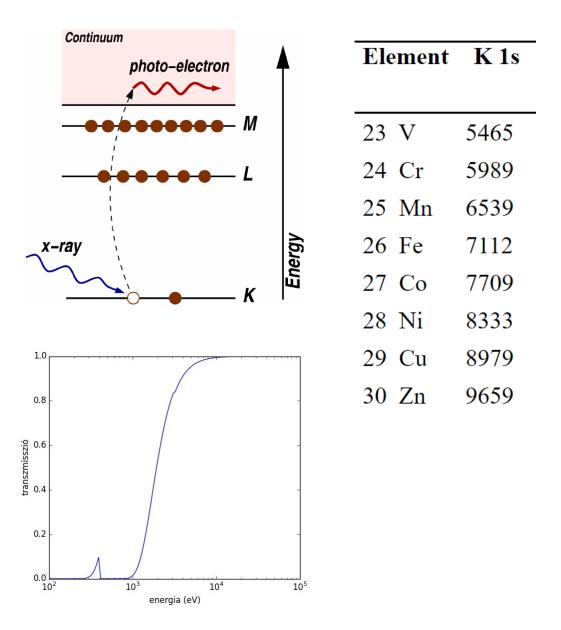
Spectroscopy with (hard) X-rays



NMR (δ)FTIR (cm⁻¹)X-ray spectroscopies (eV)FTMW (MHz)UV-Vis (nm)Mössbauer spectroscopy (mm/s)

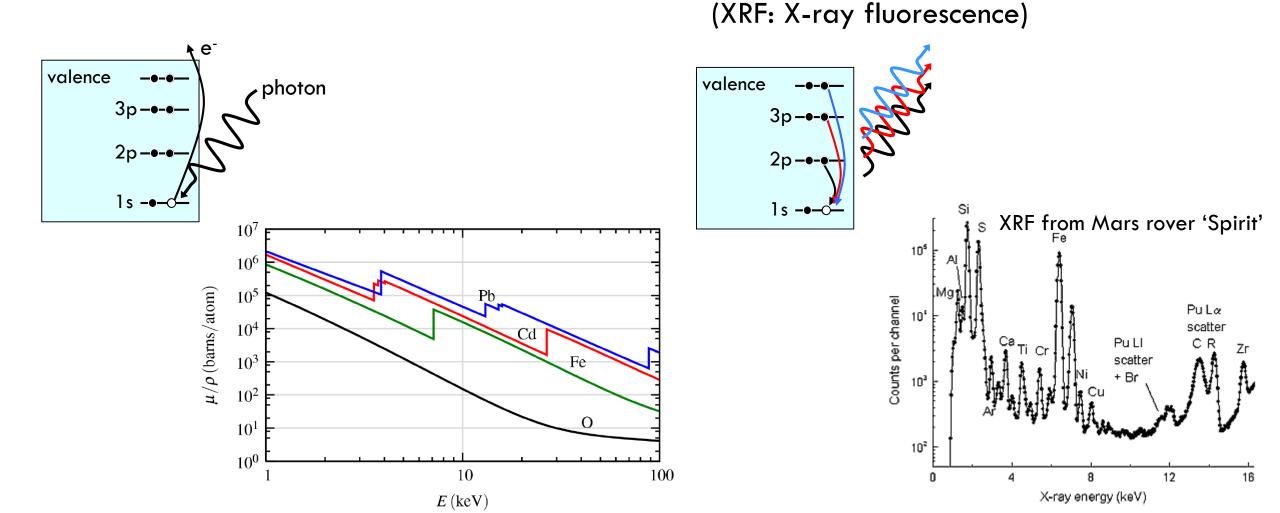
Hard X-rays

- Typically: photoexcitation of core electrons of 3d/4d metals (e.g. 1s orbital of Fe is 7112 eV, <u>http://xdb.lbl.gov/</u>)
- Elemental selectivity: energy resolution is 2 orders of magnitude better that the distance between 2 absorption edges
- High penetration depth (bulk vs. surface)
- Gas phase, liquids, solids
- Not sensitive to air \leftrightarrow UPS, XPS, soft X-rays
- But similar methodology for soft or tender X-rays which work for low Z elements like C, N, O or S, Cl, P



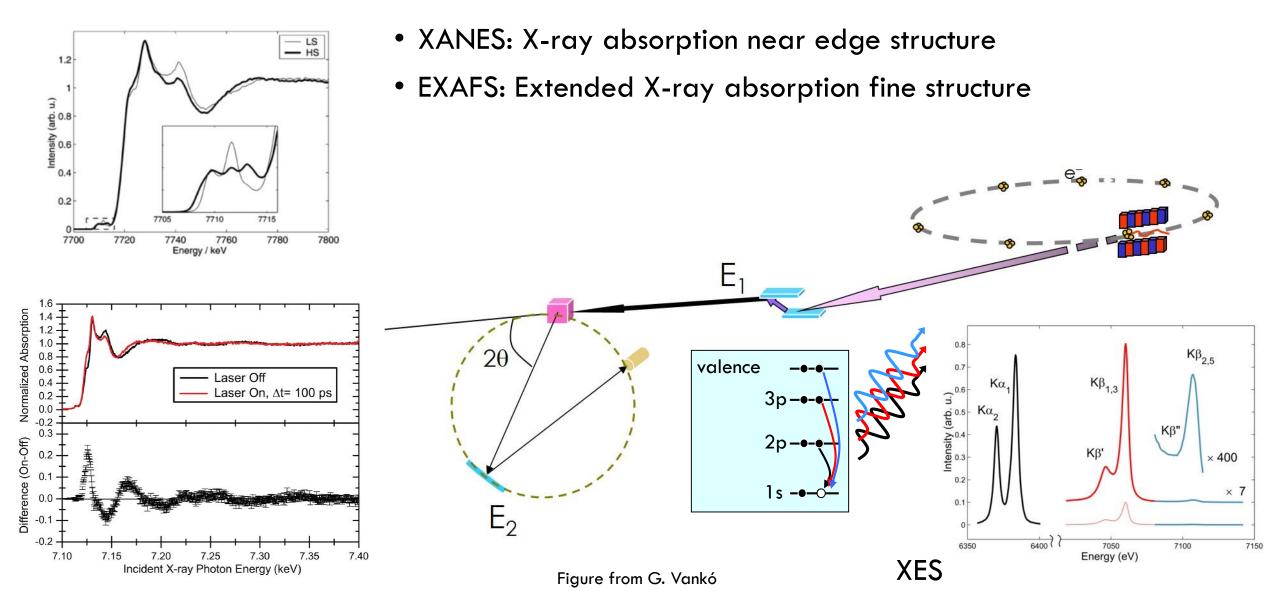
Absorption and emission of X-ray photons

XAS: X-ray absorption spectroscopy



XES: X-ray emission spectroscopy

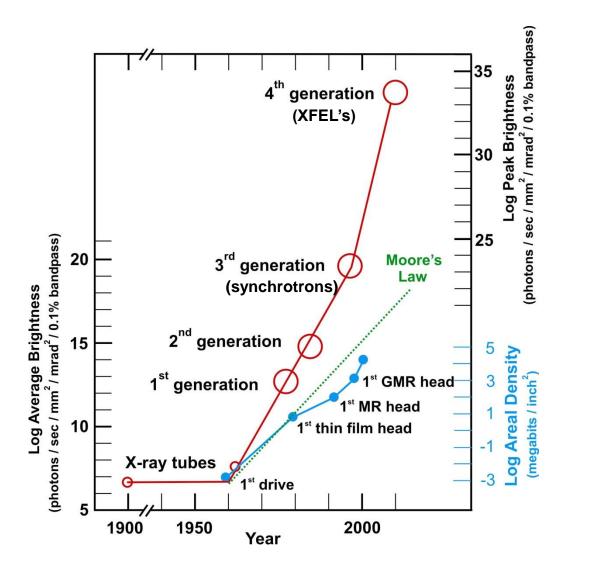
High energy resolution: reveals the fine structure

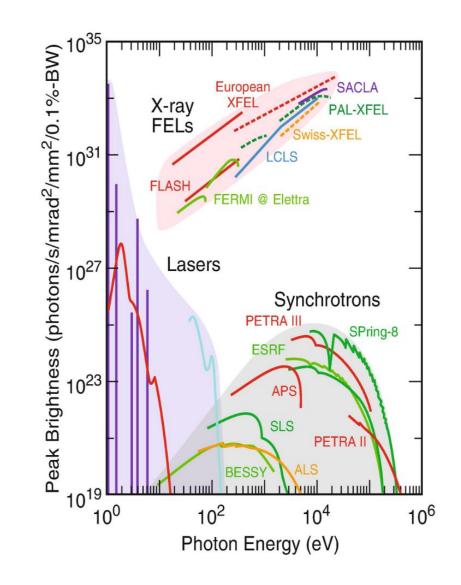


How to get these spectra?

- X-ray sources
 - Synchrotrons
 - XFEL's
 - Laboratory sources
 - Laser based sources
- Detectors
 - Ionization chambers
 - Scintillators
 - Semiconductors
- Resolving by energy (or better wavelength)
 - Internal energy resolution of the detectors
 - Bragg diffraction on single crystals
 - Laue diffraction

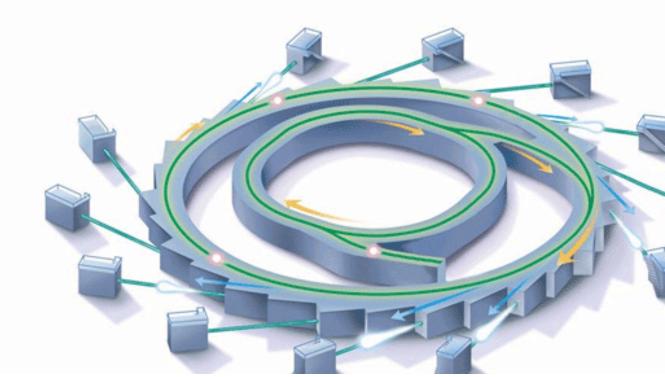
Evolution of X-ray sources





3rd gen synchrotrons (90's -)

- high brilliance: 10⁶–10¹² times more than conventional laboratory X-ray tubes
- wide energy range (from infrared to hard X-rays)
- usually linearly polarized in the ring plane
- collimated (vertical divergence: 20 μrad)
- pulsed radiation (pulse length: 10⁻¹¹–10⁻¹⁰ s, period time: 10⁻⁶–10⁻⁸ s (tunable))
- high demand and
 high operational cost ca. €100M/a
 → €10k/day/experiment



3rd gen synchrotrons (90's -)

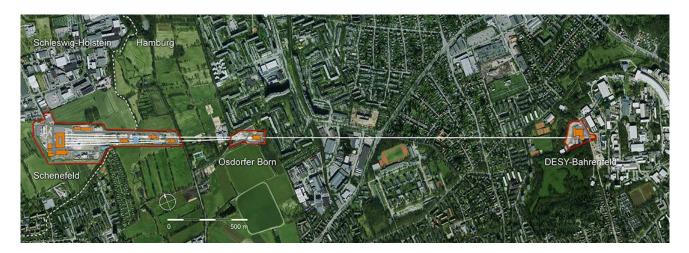
- 4 major synchrotrons (ESRF/EU, PETRA IV/DE, Spring8/Japan, APS/USA) and many smaller ones
- Beamtimes via accepted proposals (2/3 rounds per year, evaluated by international committees)
- Short, condensed beamtimes (few days)
- Local scientific and technical support
- Financial support through EU agencies for EU beamlines
- https://lightsources.org/

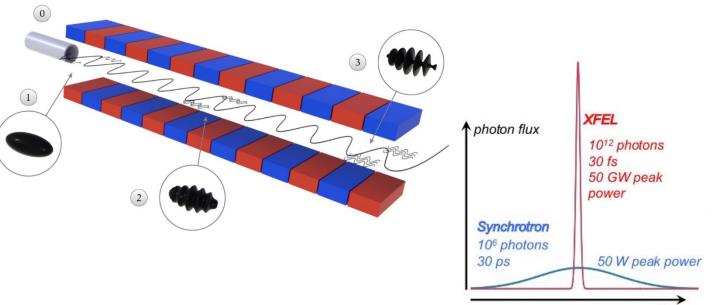




X-ray free electron lasers (10's -)

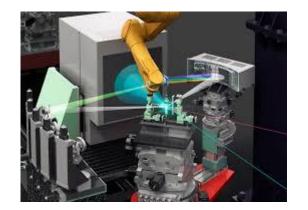
- Long linear accelerators (few km), sometimes made of superconductors
- Nearly the speed of light electrons are driven through special magnets (undulators) to emit photons
- SASE (self-amplified spontaneous emission): the coherent electron and photon bunches align each other to form ultrashort pulses
- Result: ~10 fs long extreme intense (10⁸ more peak brightness than synchrotrons) hard X-ray bunches

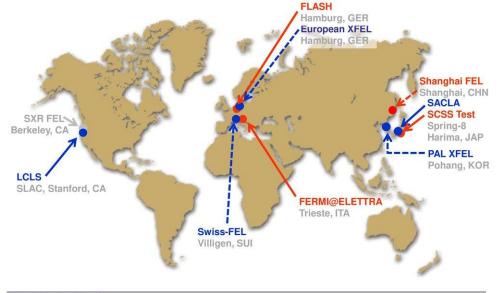




X-ray free electron lasers (10's -)

- 5 hard X-ray free electron lasers in operation (EuXFEL/EU, SwissFEL/CH, SACLA/Japan, PAL-XFEL/Korea, LCLS/USA)
- Extreme high load on the beamlines
- €1B construction cost
- €1M/day/experiment



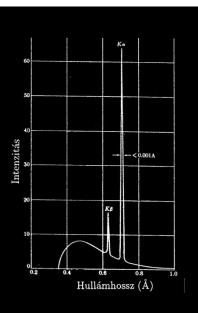


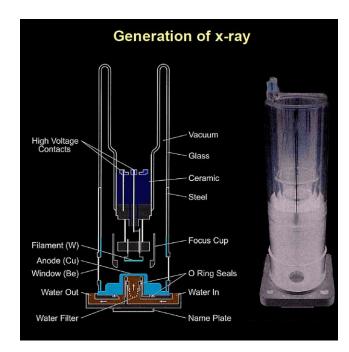
Thomas Tschentscher, European XFEL ESI 2013, CERN, 27/05/2013

Laboratory X-ray sources

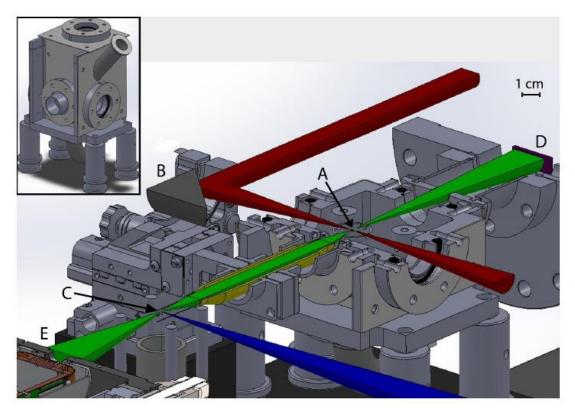
- €0.01M investment, negligible running costs
- Stable, well controlled intensity for long runs
- Can be focused to ${<}100 \mu \text{m}$
- Not polarized, no time structure (CW)
- Broad bremsstrahlung + characteristic peaks of the anode
- Intended use: diffraction, imaging, XRF







Laser plasma sources

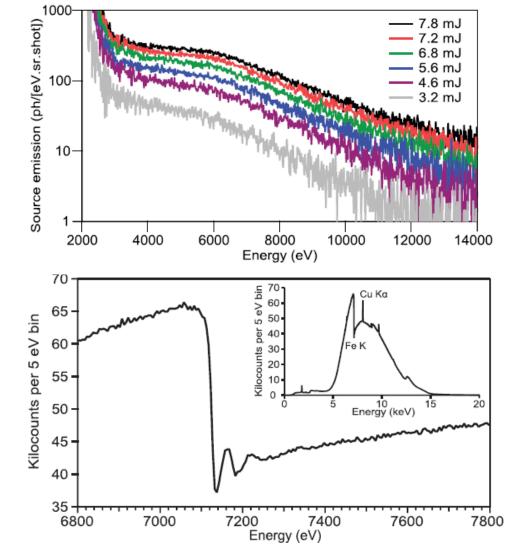


A: water jet B: mirror focusing the incident laser beam

C: sample

D: beam monitor

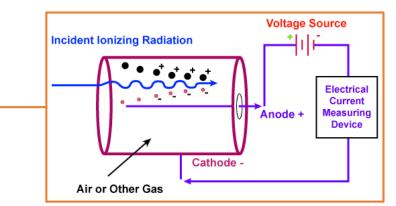
E: detector

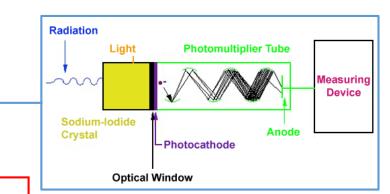


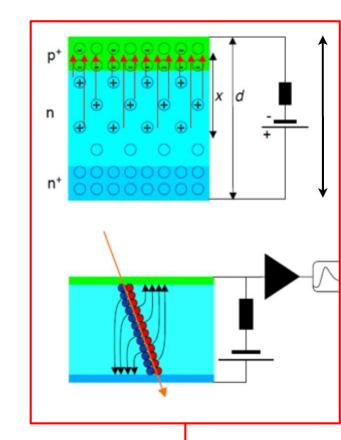
W. Fullagar, et al, Rev. Sci. Instrum. 78, 115105 (2007), Lund, Sweden L. Miaja-Avila, et al., Struct. Dyn. 2, 024301 (2015), Boulder, Colorado, USA

Detection

- Ionization chambers: -
 - Handle high photon fluxes
 - Good proportionality
- Scintillation detectors
 - Moderate energy resolution
 - Good efficiency
 - Slow
- Semiconductors
 - High energy resolution
 - Good efficiency
 - Fast (low deadtime)
 - Low background
 - Scalable can easily be used for building 2D arrays

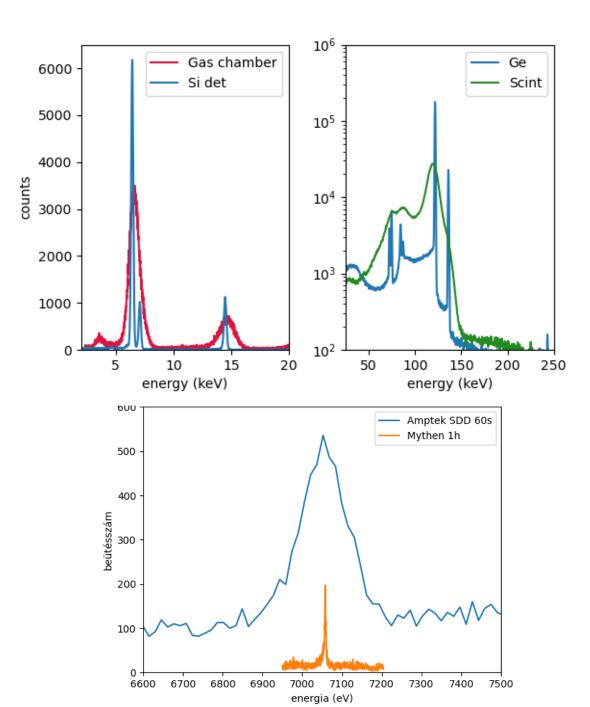






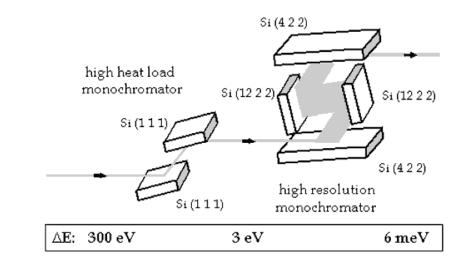
Resolving by energy

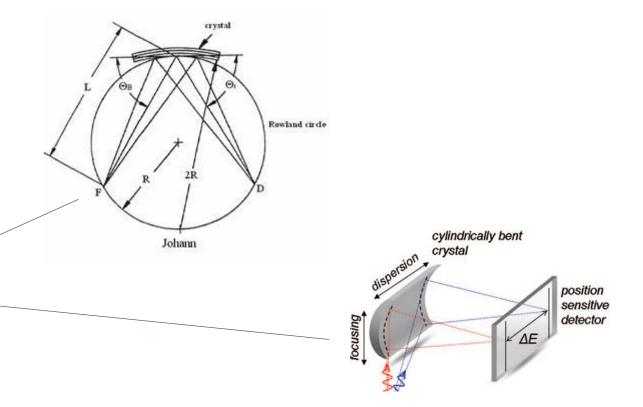
- Internal energy resolution of the detectors: discriminating by the signal height of the amplified electronic signal
- Depends strongly on the detector's working principle
- Typically: scintillators < ion chambers < semiconductors
- Best energy resolution: Silicon Drift Diodes (SDD),
 ca. 10² eV at several kV's (= 10⁻¹ E/E₀)
- But! Fine structure needs at least $10^{-3} E/E_0$



Resolving by wavelength

- Bragg diffraction or Laue diffraction
- Well know (d) diffracting single crystal analyzer
- Photons with different wavelengths will diffract at different Bragg angles
- Subsequent monochromators at high intensity synchrotron beamlines
- Smaller single crystals (Si,Ge, etc.) in spectrometers
 - Scanning mode (Rowland circle, Johann geometry with spherically bent crystal)
 - Static mode (von Hámos geometry with cylindrical crystal)





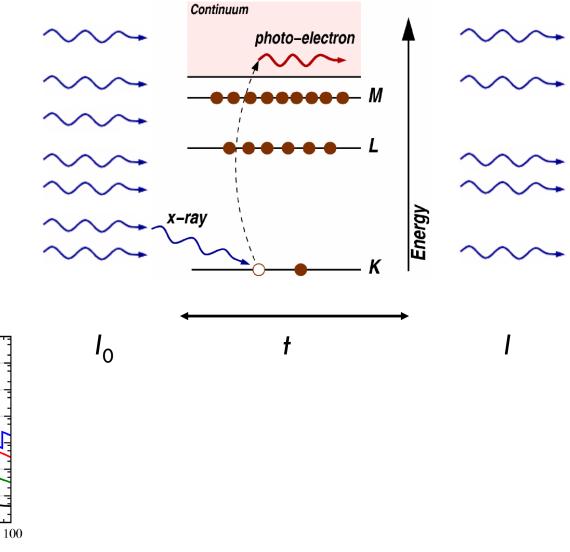
2. X-ray Absorption Spectroscopies

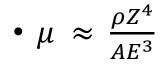
Literature:

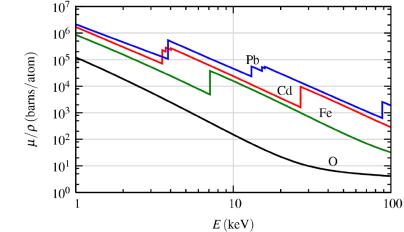
- M. Newville: Fundamentals of XAFS ,Reviews in Mineralogy & Geochemistry Vol. 78 pp. 33-74, 2014
- http://gbxafs.iit.edu/training/xafsoverview.pdf
- X-Ray Absorption Principles, Applications, Techniques of EXAFS, SEXAFS and XANES, ed. by D.
 C. Koningsberger and R. Prins, John Willey (1988)
- A. L. Ankudinov, B. Ravel, J. J. Rehr, S. D. Conradson, Phys. Rev. B 58, 7565 (1998)
- F. de Groot: High-Resolution X-ray Emission and X-ray Absorption Spectroscopy, Chem. Rev. 2001, 101, 1779-1808

The X-ray absorption effect

- Photo-electric effect (Albert Einstein, Nobel Prize in Physics, 1921)
- Beer–Lambert law: $I = I_0 e^{-\mu t}$,
- Quantity to measure: μ vs. E







The X-ray absorption coefficient

• Photoabsorption via Fermi's golden rule:

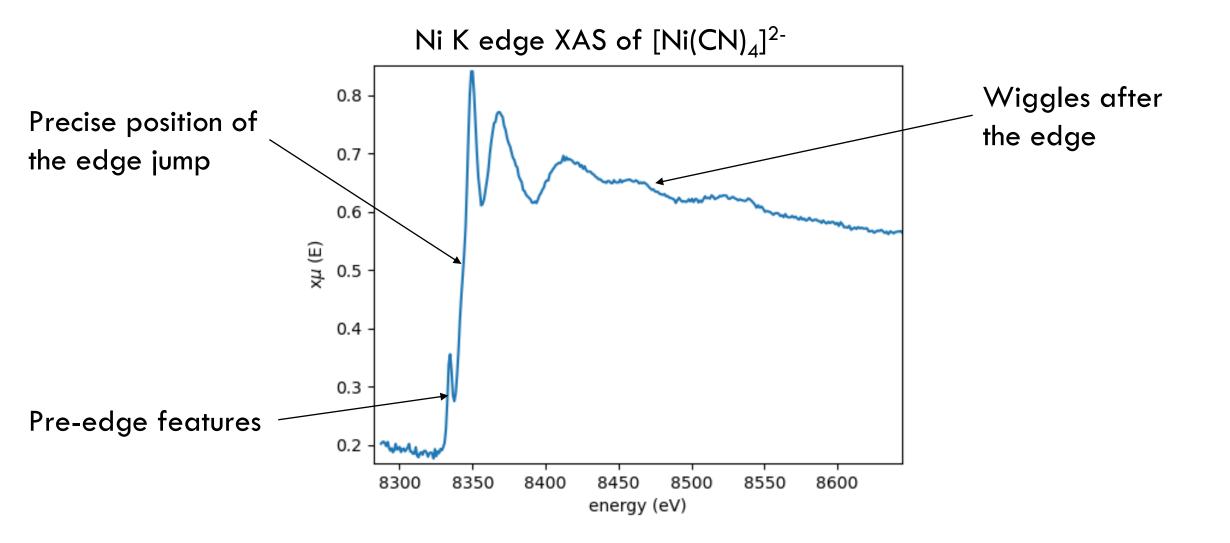
$$\mu \propto \left| \langle f \left| \hat{e} \cdot r e^{ik \cdot r} \right| i \rangle \right|^2 \delta_{E_f - E_i - \hbar \omega}$$

- Final state = initial state with a continuum electron (ε) added and a core electron removed (c)
- All inactive electrons neglected, thus the series of delta functions become the density of states (ρ):

$$u \propto |\epsilon|\hat{e} \cdot r|c\rangle|^2 \cdot \rho$$

• XAS intensity becomes a measure of the DOS

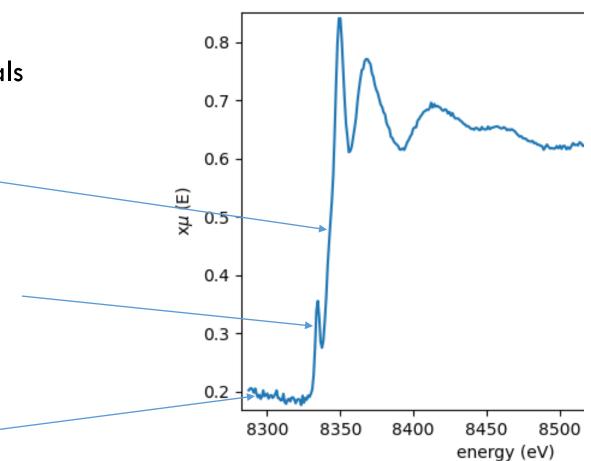
What can be seen with high energy resolution?



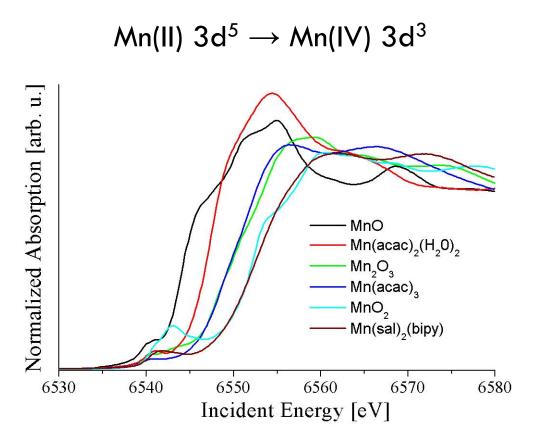
XANES – region near the absorption edge

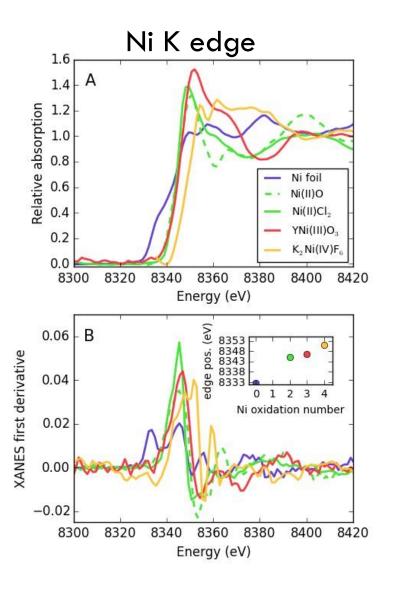
- Dipole transition matrix is ca. 137 times stronger than quadrupole, typical case is to excite 1s electron to empty 4p orbital in 3d transition metals (or 2p → 5d)
 - Edge position reflects the lowest lying empty orbitals (LUMO or conduction band)
 - ➢ Pre-edge peaks appear due to mixing the quadrupole and dipole orbitals (1s → 3d'ish)
- Absorption is small if $\hbar \omega < (E_f E_i)$

Baseline before the edge



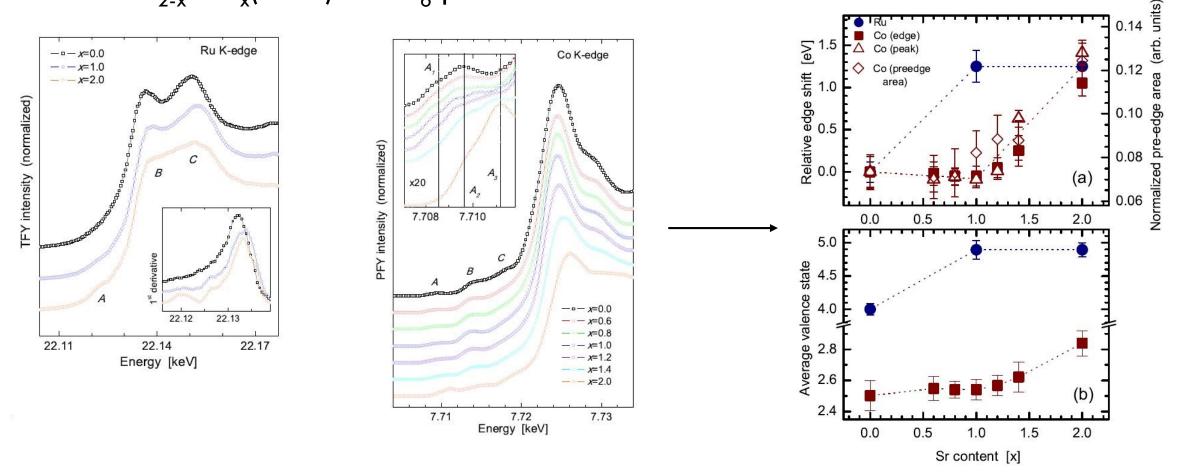
XANES – typical applications





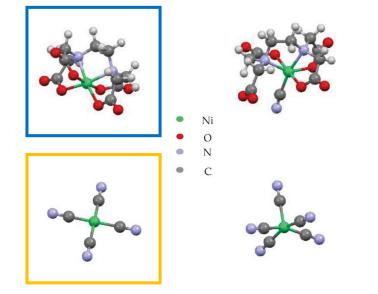
XANES – typical applications

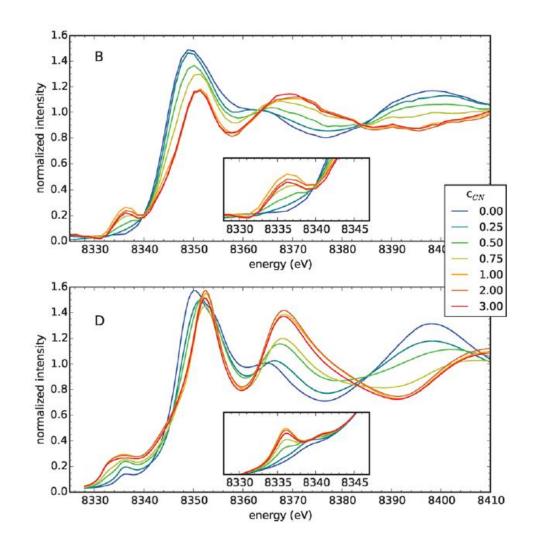
 $La^{3+}_{2-x}Sr^{2+}_{x}(CoRu)^{(6+x)+}O_{6}$ perovskite



XANES – typical applications

- Pre-edge vanishing with increasing symmetry
- From square planar to octahedral Ni²⁺



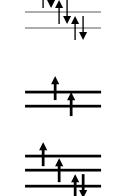


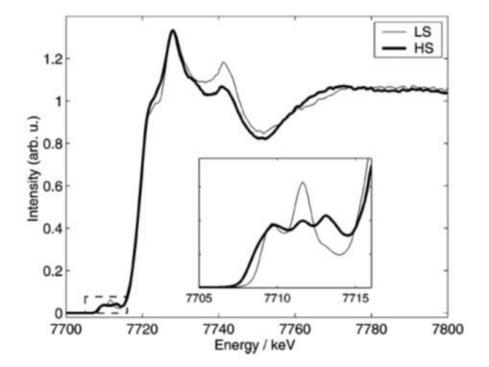
XANES – spin state dependence

- Not only the density of states, but also the spin density affects the fine structure of the XANES spectrum
- An example: Co^{3+} ion (3d⁶) in

low spin state $(t_{2a}^{6}e_{a}^{0})$



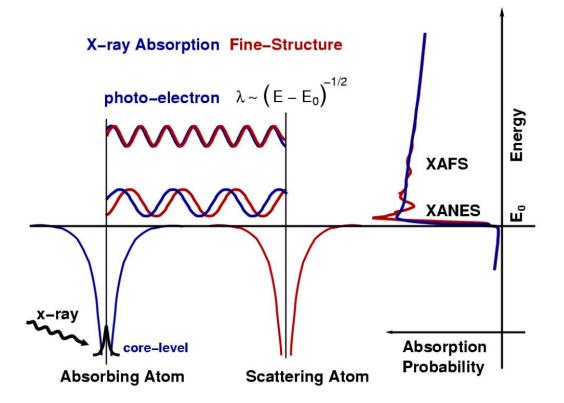




in LaCoO₃ perovskite

EXAFS – behind the absorption edge

- If the incident X-ray energy is high enough, the photoelectron will be excited to the continuum with λ
- If the excited atom has no neighbors, the absorption coefficient is not perturbed (blue line)
- The photoelectron wave can scatter back from the neighboring atoms and modify the absorption coefficient resulting in an interference pattern (red line)



The EXAFS equation

• Removing the isolated atom contribution (μ_0) and the edge jump gives the fine structure function:

$$\chi(E) = \frac{\mu(E) - \mu_0(E)}{\Delta \mu_0(E)}$$

- It is converted to wave number scale $\chi(k)$ and usually multiplies by k^2 or k^3 to account for the dumping nature of $\chi(k)$
- The final equation is:

$$\chi(k) = \sum_{i} N_i F_i(k) \frac{S_0^2}{kR_i^2} e^{\frac{-2R_i}{\lambda}} e^{-2\sigma_i^2 k^2} \sin(2kR_i + \varphi_i(k))$$

Scattering amplitude Disorder Damping

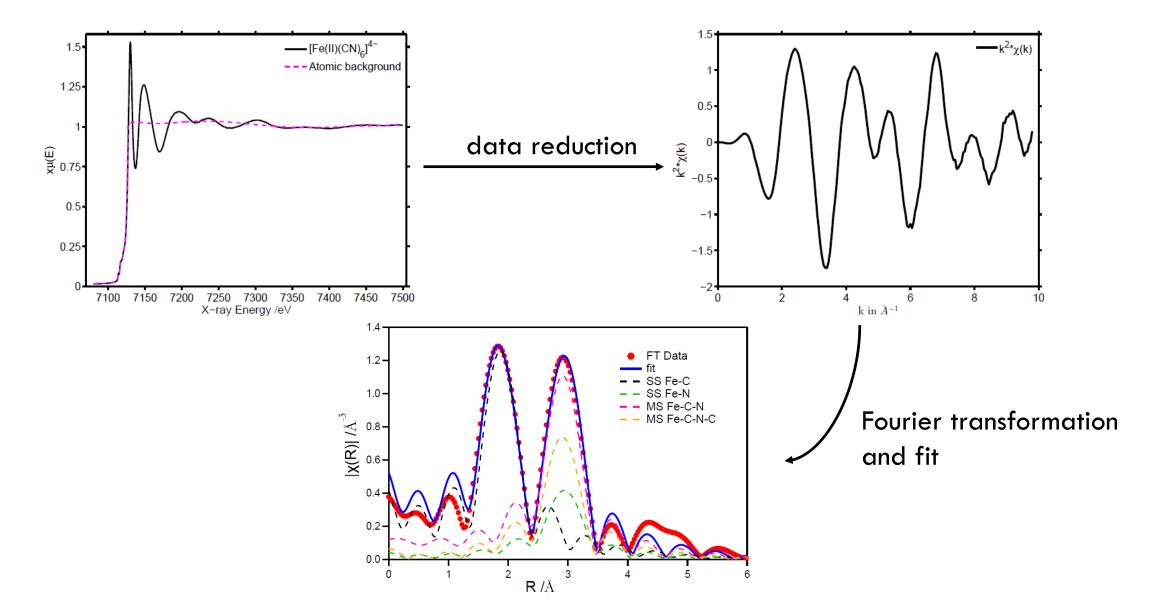
The EXAFS equation

$$\chi(k) = \sum_{i} N_i F_i(k) \frac{S_0^2}{kR_i^2} e^{\frac{-2R_i}{\lambda}} e^{-2\sigma_i^2 k^2} \sin(2kR_i + \varphi_i(k))$$

Parameters to handle:

- R_i distance between the absorber and scatterer (in case of multiple scattering, half of the total travelling path)
- N_i number of scattering atoms in the *i*th shell (=coordinantion number)
- F_i effective backscattering amplitude
- S_0^2 amplitude reduction factor
- λ mean free path of the photoelectron (the probability of returning before the core hole fills up or the photoelectron scatters inelastically)
- σ_{i}^{2} Debye-Waller factor
- $\boldsymbol{\varphi}_i$ effective scattering phase shift

EXAFS example: $[Fe^{II}(CN)_6]^{4-1}$



3. X-ray Emission Spectroscopies

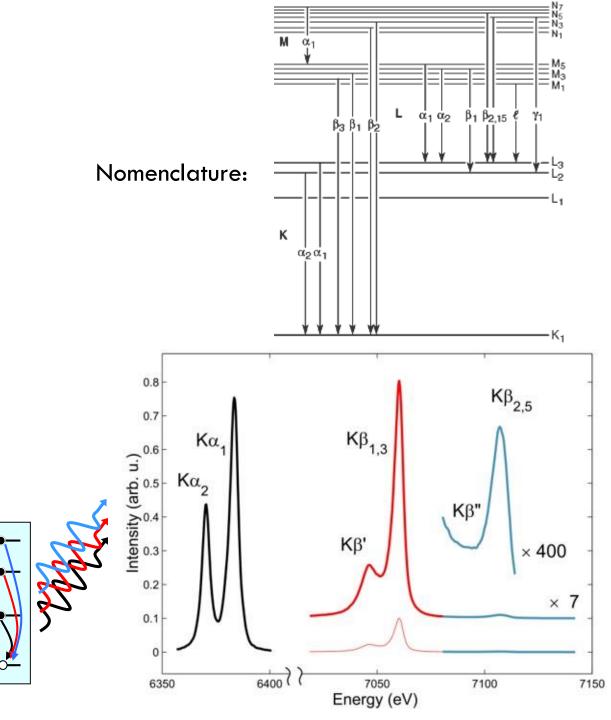
Literature:

- F. de Groot: High-Resolution X-ray Emission and X-ray Absorption Spectroscopy, Chem. Rev. 2001, 101, 1779-1808
- P. Glatzel, U. Bergmann: High resolution 1s core hole X-ray spectroscopy in 3d transition metal complexes-electronic and structural information, Coordination Chemistry Reviews 249 (2005) 65–95

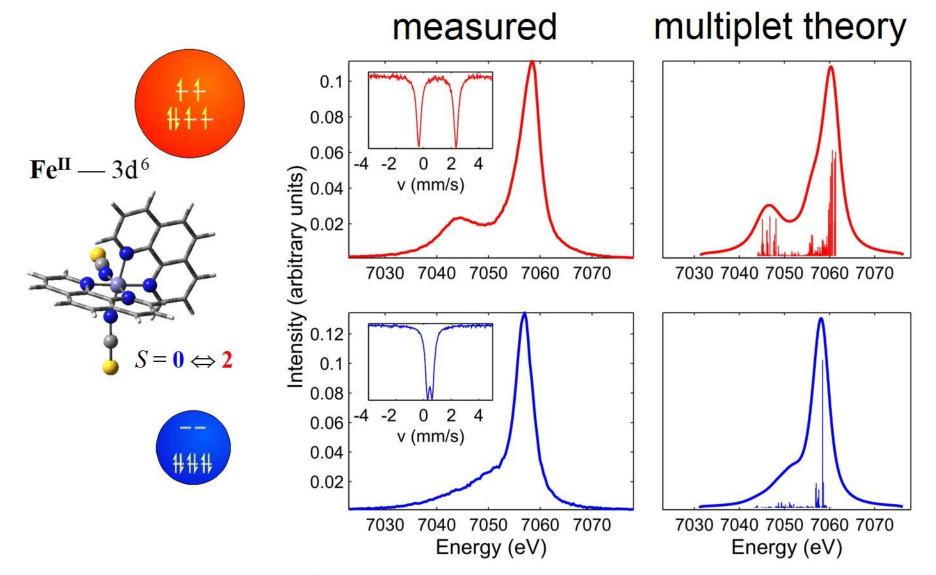
X-ray emission

- Photon-in photon-out process: after creating the core hole, an electron from outer orbitals will fill the hole, the energy difference will be (sometimes) emitted by an X-ray photon
- Core-to-core features (1s2p, 1s3p) sensitive mostly to the spin density via electron-electron exchange interaction, which is bigger on Kβ
- Spin-orbit coupling is stronger on $K\alpha_1$ and $K\alpha_2$ lines (associated with $2p_{3/2}$ and $2p_{1/2}$, respectively, typically 15-20 eV), but weaker on K β (ca. 1-2 eV)

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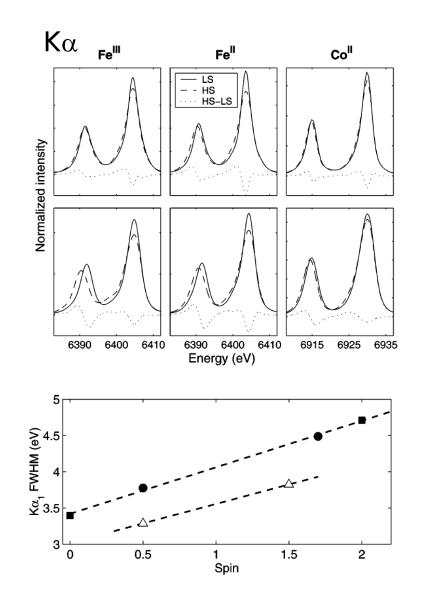
Sensitivity of ctc XES

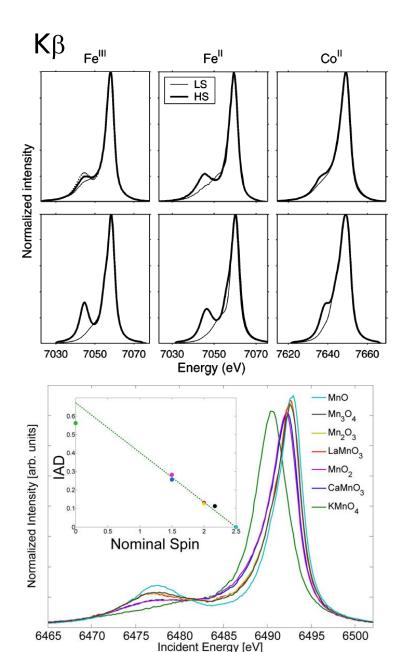


G. Vankó, F. de Groot, Phys. Rev. B 75 (2007) 177101.

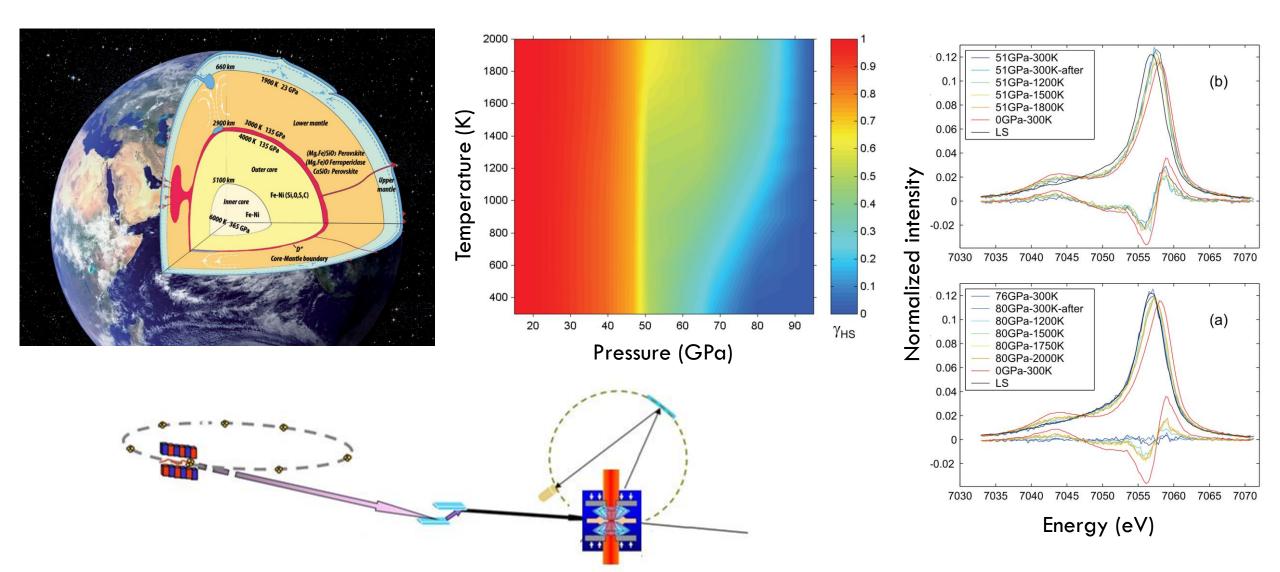
Evaluating XES spectra

- After rigorous normalization peak widths can be used to track spin state
- IAD: integrated absolute difference, footprint compared to references

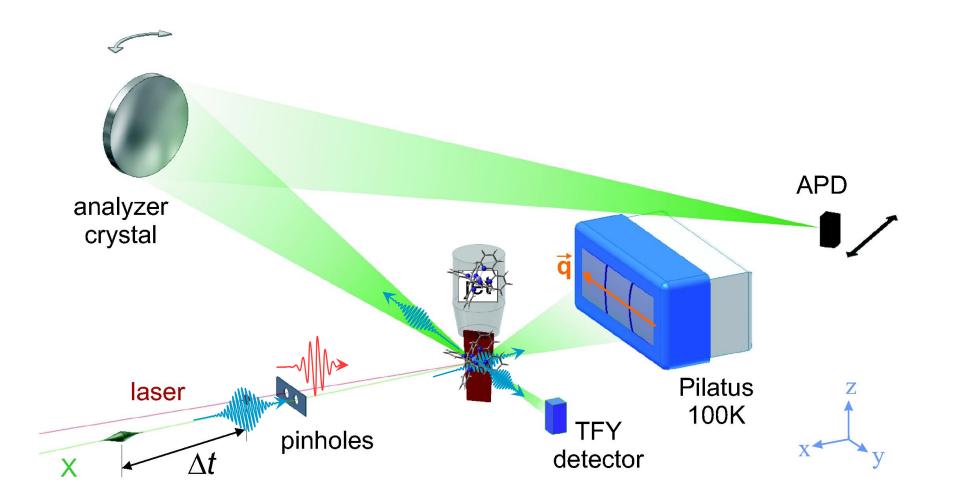




XES example: Earth's mantle

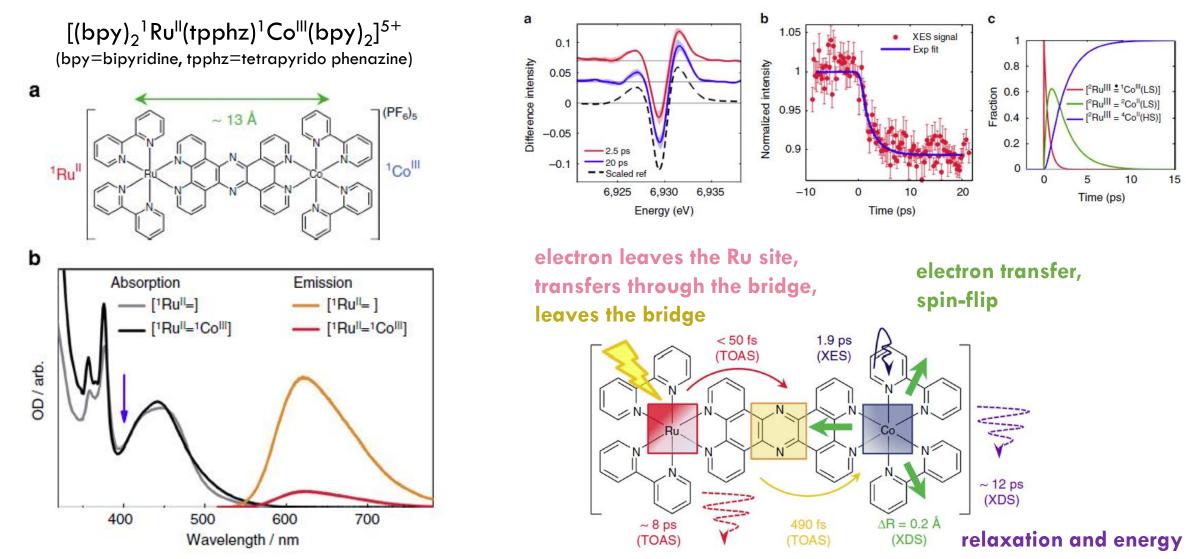


XES example: ultrafast processes



Courtesy of G. Vankó

XES example: ultrafast processes

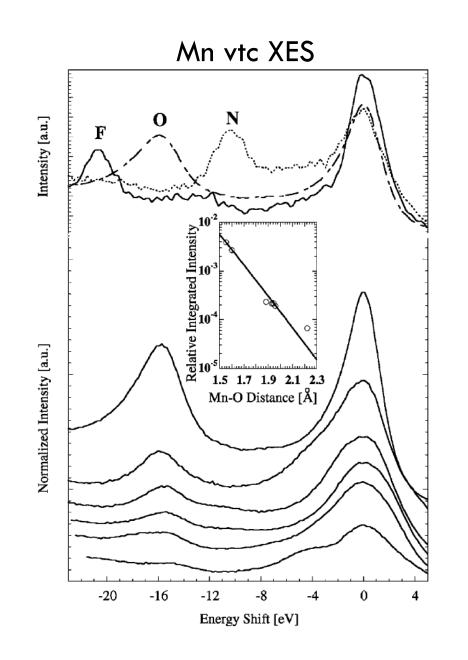


release into the solvent

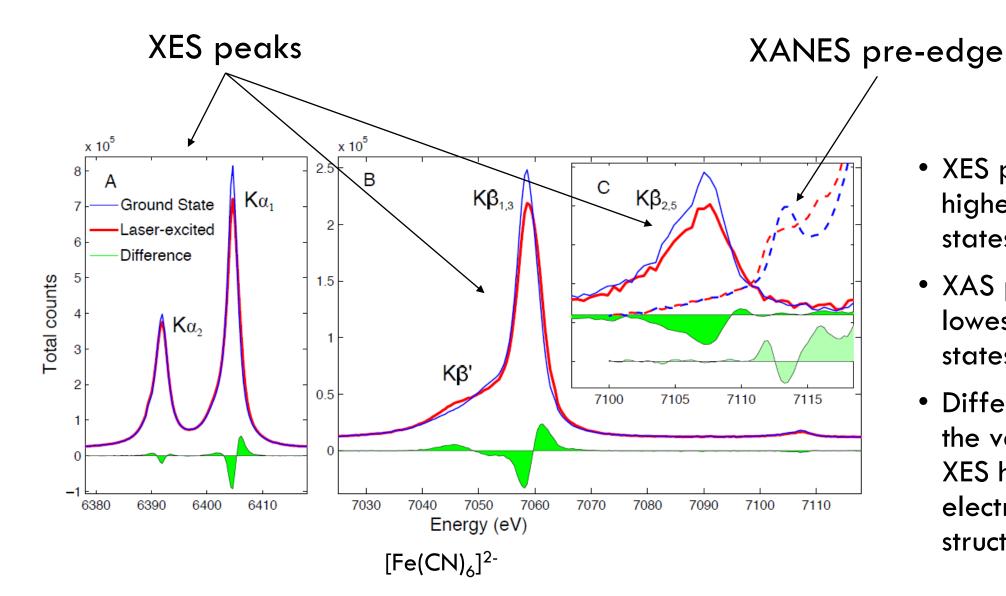
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Valence-to-core XES

- valence orbitals: chemistry!
- vtc XES reflects ligand orbitals mixed with the core ion's orbitals
- XPS has the same final state but shows also the 3d orbitals of the core ion (different selection rules!)
- vtc XES is bulk sensitive



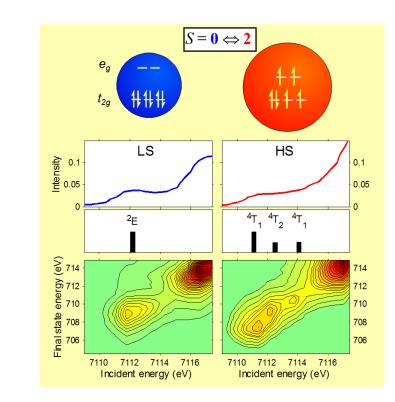
From XES to XAS



- XES probes the highest occupied states
- XAS probes the lowest unoccupied states
- Difference would be the valence gap, but XES has a different electronic core structure (core hole!)

Combining XAS and XES

- 1s2p_{3/4} resonant XES spectra
- Scanning bot the incident and fluorescent X-ray energy
- Spectral features separate better



- High Energy Resolution Fluorescent detection XANES: scanning XAS via a specific fluorescent energy
- Overcoming the 1s core hole broadening

